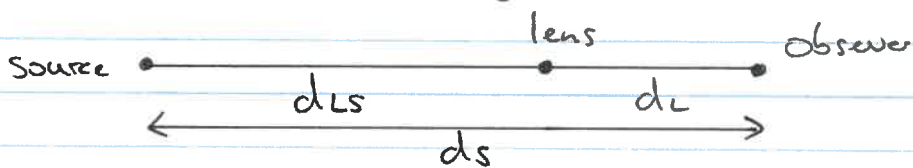


PROBLEM SET #2: SOLUTIONS

(1) Formula for angular size of Einstein ring is:

$$\theta_E = \frac{2}{c} \sqrt{\frac{GM d_{LS}}{d_L d_S}}$$

... where the distances are defined via:

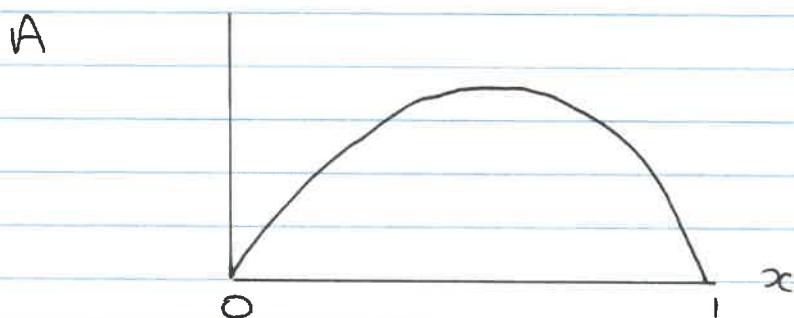


Let $d_L = x d_S$
 $d_{LS} = (1-x) d_S$.

$$\Rightarrow A = \pi d_L^2 \theta_E^2 = \frac{4\pi}{c^2} GM \frac{d_L d_{LS}}{d_S}$$

$$= \text{const} \times x(1-x)$$

3



$$A = \text{const} \times (x - x^2)$$

$$\frac{dA}{dx} = \text{const} (1 - 2x) = 0$$

3

$$\Rightarrow \underline{\underline{x_{\text{max}} = \frac{1}{2}}}$$

2

(2) Mass function for $M > M_0$:

$$\xi(M) = \xi_0 M^{-2.35}$$

Let mass of a white dwarf be M_{WD} .

Mass in remnant white dwarfs:

$$M = M_{WD} \times \text{number of white dwarfs}$$

$$= M_{WD} \int_{1M_0}^{8M_0} \xi(M) dM$$

$$= M_{WD} \xi_0 \left[\frac{M^{-1.35}}{-1.35} \right]_{1M_0}^{8M_0}$$

Total mass originally in stars:

$$M_{\text{total}} = \int_{1M_0}^{\infty} M \xi(M) dM = \xi_0 \left[\frac{M^{-0.35}}{-0.35} \right]_{1M_0}^{\infty}$$

Fraction locked up in WDs:

$$f = \frac{M}{M_{\text{total}}} = M_{WD} \times \frac{0.35}{1.35} \times \frac{\left[\frac{M^{-1.35}}{-1.35} \right]_{1M_0}^{8M_0}}{\left[\frac{M^{-0.35}}{-0.35} \right]_{1M_0}^{\infty}}$$

For $M_{WD} = 0.6 M_0$

$$f = 0.6 \times \frac{0.35}{1.35} \times \left[\frac{8^{-1.35} - 1}{-1} \right]$$

5

$$\underline{\underline{f = 0.146}}$$

(3) Mass density near the Galactic center : $\rho = \rho_0 r^{-7/4}$

(a) Total mass within 2pc is $4.5 \times 10^6 M_\odot$

$$2pc = 6.172 \times 10^{18} \text{ cm}$$

$$4.5 \times 10^6 M_\odot = 8.9505 \times 10^{39} \text{ g}$$

$$M = \int 4\pi r^2 \rho \, dr$$

$$= \int 4\pi \rho_0 r^{1/4} \, dr$$

$$= 4\pi \rho_0 \left[\frac{r^{5/4}}{5/4} \right]_0^{2pc}$$

$$\Rightarrow \rho_0 = 2.89 \times 10^{15} \text{ in c.g.s. units.}$$

$$= 2.89 \times 10^{15} \text{ g cm}^{-5/4}$$

3

2

(b) This is just $v = \sqrt{\frac{GM_{BH}}{r}}$

(c) Formula for dV/dt from the lectures is :

$$\frac{dV}{dt} = - \frac{4\pi G^2 M}{V^2} nM \ln \Lambda$$

where M is the massive object, moving with velocity v with respect to less massive stars of mass m and number density n.

$$\ln \Lambda = \ln \left(\frac{b_{max}}{b_{min}} \right) \dots \text{limits on impact parameter } b.$$

\therefore for the dynamical friction we have a scaling:

$$t_{\text{friction}} = \frac{V}{|dV/dt|} \propto \frac{V^3}{nM}$$

$$\left. \begin{array}{l} nM = \rho \propto r^{-3/4} \\ V \propto r^{-1/2} \end{array} \right\} t_{\text{friction}} \propto \frac{r^{-3/2}}{r^{-3/4}} \propto r^{1/4}$$

Now substitute numbers:

$$t_{\text{friction}} = \frac{V^3}{4\pi G^2 M \rho \ln \Lambda}$$

$$M = 10 M_{\odot} = 1.989 \times 10^{34} \text{ g}$$

$$\begin{aligned} \rho &= \rho_0 (1 \text{ pc})^{-3/4} \\ &= 1.272 \times 10^{-17} \text{ g cm}^{-3} \end{aligned}$$

$$V = 1.39 \times 10^7 \text{ cm s}^{-1}$$

$$\ln \Lambda = 10.$$

$$\Rightarrow \underline{t_{\text{friction}} \approx 6.0 \times 10^8 \text{ yr.}}$$

PROBLEM SET # 3: SOLUTIONS

(1)

$$\oint_S \underline{g} \cdot d\underline{A} = -4\pi G M_{in}$$

For a spherically symmetric mass distribution $\underline{g} = (-g_r, 0, 0)$ in polar co-ordinates. We have:

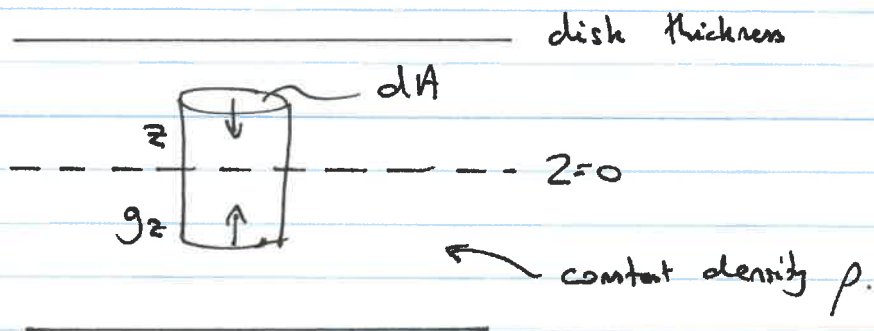
$$g_r \cdot 4\pi r^2 = 4\pi G M_{in}$$

$$g_r = \frac{G M_{in}}{r^2}$$

(5)

... as expected.

(2)



Consider a cylinder extending from $\pm z$ with cap area dA . By symmetry, $\underline{g} = 0$ except on end caps, where it equals g_z .

Applying formula above g_z (taking positive downwards):

$$2 g_z dA = 4\pi G \cdot 2z dA \rho.$$

$$\underline{\underline{g_z = 4\pi G \rho z}}$$

(1)

(3)

Since g_z is directed downwards, equation of motion is:

$$\frac{d^2 z}{dt^2} = -4\pi G\rho z$$

①

$$\frac{d^2 z}{dt^2} + k z = 0 \quad \text{with} \quad k = 4\pi G\rho.$$

②

General solⁿ will be sum of sin and cos terms.

e.g. trial solⁿ

$$z = \sin(\omega t)$$

$$\dot{z} = \omega \cos(\omega t)$$

$$\ddot{z} = -\omega^2 \sin(\omega t)$$

$$\Rightarrow -\omega^2 \sin(\omega t) + k \sin(\omega t) = 0$$

$$\omega = k^{1/2}.$$

③ ④

$$\text{Solⁿ is } z = A \cdot \sin \left[(4\pi G\rho)^{1/2} t \right] + B \cos \left[(4\pi G\rho)^{1/2} t \right]$$

For one oscillation: $\omega t = 2\pi$

$$\Rightarrow P = \frac{2\pi}{(4\pi G\rho)^{1/2}} = \left(\frac{\pi}{G\rho} \right)^{1/2}$$

$$\text{If } \rho = 0.1 M_{\odot} \text{ pc}^{-3} = 6.77 \times 10^{-24} \text{ g cm}^{-3}$$

$$\Rightarrow P = 2.6 \times 10^{15} \text{ s}$$

②

④

$$= \underline{\underline{8.35 \times 10^7 \text{ yr}}}$$

PROBLEM SET #4 - SOLUTIONS

(1) The Eddington limit is:

$$L_{\text{Edd}} = \frac{4\pi G M_p}{\sigma_e} M$$

... where σ_e is the Thomson cross-section for electron-photon scattering.

$$\text{If } L = \eta \dot{M} c^2 \text{ with } \eta = \text{const.}$$

$$\dot{M} = \frac{1}{\eta c^2} \frac{4\pi G M_p}{\sigma_e} M$$

$$\text{i.e. } \frac{dM}{dt} = kM \quad \text{with } k = \frac{4\pi G M_p}{\eta c \sigma_e}$$

$$\text{Solving this eqn: } \int_{M_0}^M \frac{dM'}{M'} = \int_0^t k dt'$$

$$\ln \left[\frac{M}{M_0} \right] = kt$$

$$M = M_0 e^{kt} = M_0 e^{t/\tau}$$

$$\dots \text{ where } \tau = k^{-1} = \frac{\eta c \sigma_e}{4\pi G M_p}$$

$$\text{Numerically, for } \eta = 0.1, \quad \tau = 1.4 \times 10^{15} \text{ s} = \underline{\underline{4.5 \times 10^7 \text{ yr}}}$$

Note: strictly need to account for fact that radiated energy does not add to rest mass of black hole. In expression for τ this means $\eta \rightarrow \eta/(1-\eta)$... but this correction is not too important.

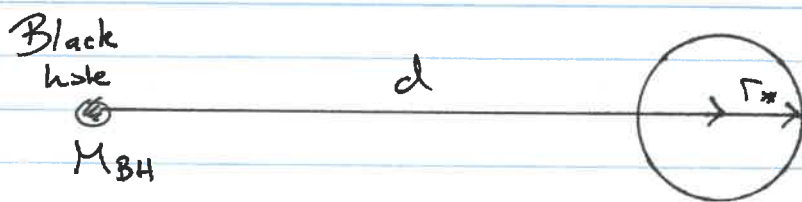
Time to grow from $10 M_{\odot} \rightarrow 10^9 M_{\odot}$ by Eddington limited accⁿ is then:

$$\ln \left[\frac{10^9}{10} \right] \times \tau = 18.4 \tau$$

$$= 8.3 \times 10^8 \text{ yr.}$$

\therefore There is not enough time available

(2)



Tidal force:
per unit mass
(i.e. tidal accⁿ)

$$F_{\text{tidal}} = \frac{GM_{\text{BH}}}{d^2} - \frac{GM_{\text{BH}}}{(d+r_*)^2}$$

$$= \frac{GM_{\text{BH}}}{d^2} \left[1 - \left(1 + \frac{r_*}{d}\right)^{-2} \right]$$

for $r_* \ll d$ this can be approximated as $1 - \frac{2r_*}{d} + \dots$

$$\Rightarrow F_{\text{tidal}} = \frac{2GM_{\text{BH}}}{d^3} \Gamma_*$$

Self-gravity force of the star:

$$F_{\text{sg}} = \frac{GM_*}{r_*^2} \quad (\text{again force per unit mass})$$

Equating, the critical distance is:

$$d = \left(\frac{2M_{\text{BH}}}{M_*} \right)^{\frac{1}{3}} \Gamma_*$$

... with disruption @ smaller radii. As quoted ignoring the $2^{\frac{1}{3}}$ numerical constant.

For the Galactic Center: $M_{\text{BH}} \approx 4.5 \times 10^6 M_{\odot}$

$$d \approx \left(\frac{4.5 \times 10^6 M_{\odot}}{M_{\odot}} \right)^{\frac{1}{3}} \Gamma_{\odot}$$

\uparrow
 $7 \times 10^{10} \text{ cm}$

$$\Rightarrow \text{tidal radius } d \approx 1.15 \times 10^3 \text{ cm.}$$

$$\text{Schwarzschild radius } R_s = \frac{2GM_{\text{BH}}}{c^2} = 1.3 \times 10^{12} \text{ cm.}$$

$$\Rightarrow \text{Tidal radius is at } \approx \underline{\underline{8.6 \text{ Schwarzschild radii}}}$$