

Appendix D

Glossary of Symbols

A prime (') generally denotes interaction-picture operators.

V_0	rf voltage amplitude
Ω_T	rf drive frequency
d_0	characteristic trap dimension
Q	ion charge
m	ion mass
z_{tot}	ion position, including secular motion and micromotion
z	secular motion
z_μ	micromotion
$x_{\mu 0, j}$	micromotion amplitude of j^{th} ion in x -direction (2-ion entanglement)
ω_m	secular frequency of mode m
ω_z	z secular frequency $\approx \frac{2\sqrt{2}QV_0}{md_0^2\Omega_T}$
ω_r	radial secular frequency (spherical quadrupole trap) $\omega_r = \omega_z/2$
\bar{D}_z	pseudopotential well depth (z -direction) $\approx \frac{QV_0^2}{md_0^2\Omega_T^2}$
U_0	static quadrupole voltage on trap electrodes
q	Mathieu q-parameter
a	Mathieu a-parameter

β	$\frac{2\omega_m}{\Omega_T} \approx \sqrt{a + \frac{g^2}{2}}$
	<i>also</i> displacement parameter due to noisy electric fields
	<i>also</i> qubit coefficient
\hat{a}	harmonic oscillator lowering operator = $\sqrt{\frac{m\omega_z}{2\hbar}}(\hat{z} + \frac{i}{m\omega_z}\hat{p})$
\hat{a}^\dagger	harmonic oscillator raising operator = $\sqrt{\frac{m\omega_z}{2\hbar}}(\hat{z} - \frac{i}{m\omega_z}\hat{p})$
\hat{n}	harmonic oscillator number operator = $\hat{a}^\dagger\hat{a}$
<i>H.C.</i>	Hermitian conjugate
z_0	ground state wave packet spread = $\sqrt{\frac{\hbar}{2m\omega_z}}$
\mathcal{Q}	resonator quality factor
Z_0	resonator characteristic impedance
α	atomic polarizability of neutral background gas atoms
	<i>also</i> coherent state amplitude
	<i>also</i> qubit coefficient
μ_{red}	reduced mass for ion/neutral collisions
ν	relative velocity in ion/neutral collisions
b	impact parameter for ion/neutral collisions
b_{crit}	critical impact parameter for spiralling ion/neutral collisions
$k_{Langevin}$	Langevin rate constant
$\gamma_{Langevin}$	reaction rate due to background neutral gas particles/spiralling collisions
γ_{el}	background gas reaction rate based on total elastic cross section
ρ	background gas density
$\hat{\rho}, \rho_{n,m}$	density matrix, matrix elements
σ_{el}	total elastic collision cross section
$\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z$	unit vectors in the x -, y -, and z - directions

e	electron charge
$ \downarrow\rangle$	${}^9Be^+$, $2s\ 2S_{1/2}$, $ F=2, m_F=-2\rangle$ state (sometimes $ F=2, m_F=2\rangle$)
$ \uparrow\rangle$	${}^9Be^+$, $2s\ 2S_{1/2}$, $ F=1, m_F=-1\rangle$ state (sometimes $ F=1, m_F=1\rangle$)
$ v\rangle$	“virtual level” for Raman transitions: ${}^9Be^+$, $2s\ 2S_{1/2}$ level
γ	excited state ($2p$) linewidth ≈ 19.4 MHz for ${}^9Be^+$ <i>also</i> Berry’s (geometrical) phase
δ_{FS}	${}^9Be^+$ fine-structure splitting ≈ 197 GHz
τ	excited state lifetime ($= 1/2\pi\gamma$)
$\hat{\sigma}_1, \hat{\sigma}_2, \hat{\sigma}_3$	Pauli matrices (one also sees the notation $\hat{\sigma}_x, \hat{\sigma}_y, \hat{\sigma}_z$)
$\hat{S}_1, \hat{S}_2, \hat{S}_3$	atomic spin operators $= \frac{1}{2}\hbar\hat{\sigma}_j$
\hat{S}_+, \hat{S}_-	$\hat{S}_\pm = \frac{1}{2}(\hat{S}_1 \pm \hat{S}_2)$
$\hat{\mu}$	electric dipole moment of the atom $= e\hat{\mathbf{r}}_{el}$
\mathbf{k}, k	laser wave vector, magnitude
$\Delta\mathbf{k}$	laser wave vector difference (Raman transitions)
ω_L	laser frequency
$\Delta\omega_L$	laser frequency difference
ϕ	laser phase <i>also</i> general phase
ω_0	carrier resonance frequency (${}^9Be^+$ hyperfine splitting plus Zeeman shift difference)
δ	laser detuning from resonance
Δ_R	Raman laser detunings from virtual level v

Ω	carrier Rabi frequency
$\Omega_{n,m}$	sideband Rabi frequency
Ω_d	effective displacement amplitude for applied electric fields
\hat{g}_i	single-photon coupling operators (Raman transitions) $= E_i e^{-i\phi_i} \boldsymbol{\epsilon}_i \cdot \hat{\mathbf{x}}/2\hbar$
g_1	single-photon coupling strength (Raman transitions) $= \langle \downarrow \hat{g}_1 v \rangle$
g_2	single-photon coupling strength (Raman transitions) $= \langle \uparrow \hat{g}_2 v \rangle$
g	general Rabi frequency (e.g. if $g_1 = g_2 = g$) <i>also</i> red sideband coupling strength for engineered $T = 0$ reservoir $= i\eta\hbar\Omega e^{i\phi}$
η	Lamb-Dicke parameter ($= kz_0$, 1D)
\bar{n}	average phonon number
ζ	overall detection efficiency <i>also</i> dimensionless time in Mathieu equation
t_{DG}	length of Detection Gate
t_{pr}	length of probe pulse
T_R	length of free evolution time in Ramsey experiment
I	laser intensity
I_S	saturation intensity (≈ 85 mW/m ² for ${}^9\text{Be}^+$)
\bar{m}	avg. number of photons detected
$\hat{D}(\alpha)$	displacement operator $= e^{(\alpha\hat{a}^\dagger - \alpha^*\hat{a})}$
L_n^m	associated Laguerre polynomial
P_\downarrow	probability that the ion is in $ \downarrow\rangle$

$C_n, C_{\downarrow,n}, C_{\uparrow,n}$	motional state probability amplitude
$C_i^{(k)}(t)$	k^{th} -order perturbative expansion of $C_i(t)$
P_n	prob. that the ion is in motional Fock state $ n\rangle$
r	ratio of red to blue sideband sizes <i>also</i> squeeze amplitude in \hat{K}_1, \hat{K}_2 formalism
p_r	recoil momentum = $\hbar k$
E_r	recoil energy = $\frac{p_r^2}{2m}$
\mathcal{B}	background counts
$S_E(\omega)$	electric field noise spectral density (V/m ²)
Γ_0	rate at which ion is heated out of ground state of motion
k_B	Boltzmann's constant $\approx 1.38 \times 10^{-23}$ J/K
$g(t)$	time-dependent laser coupling strength (“walking standing wave” generation of coherent states) $= \langle \downarrow \mathbf{e}_x \cdot \mathbf{r}_{el} \frac{eE}{2\hbar} \cos[\frac{1}{2}(\Delta kz - \delta t + \delta\phi)] v \rangle$
$\hat{S}(\epsilon)$	squeeze operator = $\exp[\frac{1}{2}(\epsilon^* \hat{a}^2 - \epsilon(\hat{a}^\dagger)^2)]$
ϵ	squeeze parameter $\epsilon = R e^{2i\phi}$ <i>also</i> energy
$\hat{\epsilon}$	laser polarization vector
Ω_{\downarrow}	Rabi frequency for $ \downarrow\rangle$ (including laser polarization/matrix elements)
Ω_{\uparrow}	Rabi frequency for $ \uparrow\rangle$ (including laser polarization/matrix elements)
$W(\alpha) \equiv W(z, ip)$	motional Wigner function
$Q_n(\alpha) \equiv \mathcal{C}_n(\alpha) ^2$	Fock state probabilities for displaced motional state (density matrix reconstruction) = $\langle n \hat{D}^\dagger(\alpha) \hat{\rho} \hat{D}(\alpha) n \rangle$

$ \psi_B^\pm\rangle$	Bell states = $\frac{1}{\sqrt{2}}(\downarrow, \uparrow\rangle \pm \uparrow, \downarrow\rangle)$
$ \psi_e(\phi)\rangle$	approximate Bell state = $\left(\frac{3}{5} \downarrow, \uparrow\rangle - e^{i\phi\frac{4}{5}} \uparrow, \downarrow\rangle\right) n=0\rangle$
Ω_i	Rabi frequency of ion i (Bell state production)
ε	detuning in “Sørensen and Mølmer” entangled-state production scheme <i>also</i> small error in Berry’s phase squeeze operators <i>also</i> error
\oplus	Exclusive-OR operator (addition modulo 2) <i>also</i> Hilbert space direct sum
$ \mathcal{R}\rangle$	state of qubit register
\mathcal{R}	parameters upon which Berry’s phase depends <i>also</i> reservoir
$\hat{\mathcal{H}}, \hat{\mathcal{H}}^{(N)}$	Hadamard transform
$ aux\rangle$	auxilliary level for Controlled-NOT logic gate $= 2s^2 S_{1/2}, F=2, m_F=0\rangle$
\hat{s}_i	system operator (reservoir interactions)
$\hat{\Gamma}_i$	reservoir operator of mode i (reservoir interactions)
$\hat{\chi}$	complete density matrix of system + reservoir
\hat{R}_0	initial reservoir density matrix
\mathcal{M}	superoperator
$\mathbb{1}$	identity operator
$\hat{b}_k, \hat{b}_k^\dagger$	lowering and raising operators for reservoir (harmonic oscillator) mode k
\bar{N}	average reservoir occupation number (for given mode or frequency)
κ	coupling strength (to reservoir mode)

\mathcal{V}	rms voltage noise
Γ	effective decay rate ($T = 0$ reservoir engineering)
Ω_p	Rabi frequency of Red Doppler beam ($T = 0$ reservoir engineering)
ξ	dynamical phase
$\hat{U}(t)$	(Ch. 9) unitary operator which changes quantum state to produce Berry's phase
$\hat{G}(t)$	(Ch. 9) operator which generates $\hat{U}(t)$
$\hat{S}(t)$	spin operator for spin- $\frac{1}{2}$ Berry's phase
\mathbf{f}_i	i^{th} frame field unit vector
$\omega_{i,j}$	differential "one form" which measures rotation of frame field
\hat{K}_1	generator of squeezes along z $= -\frac{i}{4}[\hat{a}^2 - (\hat{a}^\dagger)^2]$
\hat{K}_2	generator of squeezes along $z + ip$ $= \frac{1}{4}[\hat{a}^2 + (\hat{a}^\dagger)^2]$
$\hat{J}_3 \equiv \hat{K}_3$	generator of rotations in (z, ip) -plane $= -\frac{1}{2}(\hat{n} + \frac{1}{2})$
$\hat{K}(t)$	state operator for SU(1,1) Berry's phase
$\tilde{\times}$	Minkowski cross product (see Eq. (9.44))
$\tilde{\cdot}$	Minkowski dot product (see Eq. (9.45))
$\tilde{\varepsilon}_{ijk}$	$\doteq (-1)^{\delta_{k,3}} \varepsilon_{ijk}$