

Macroscopic Quantum Tunneling and Entangled States in Bose-Einstein Condensates



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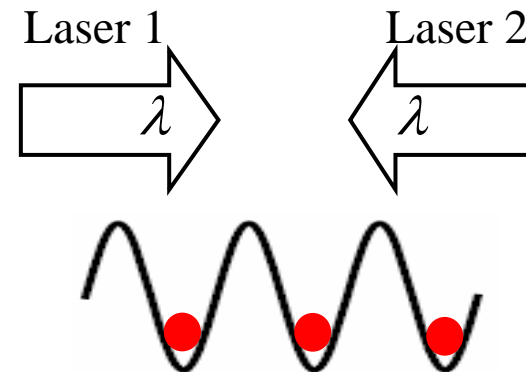
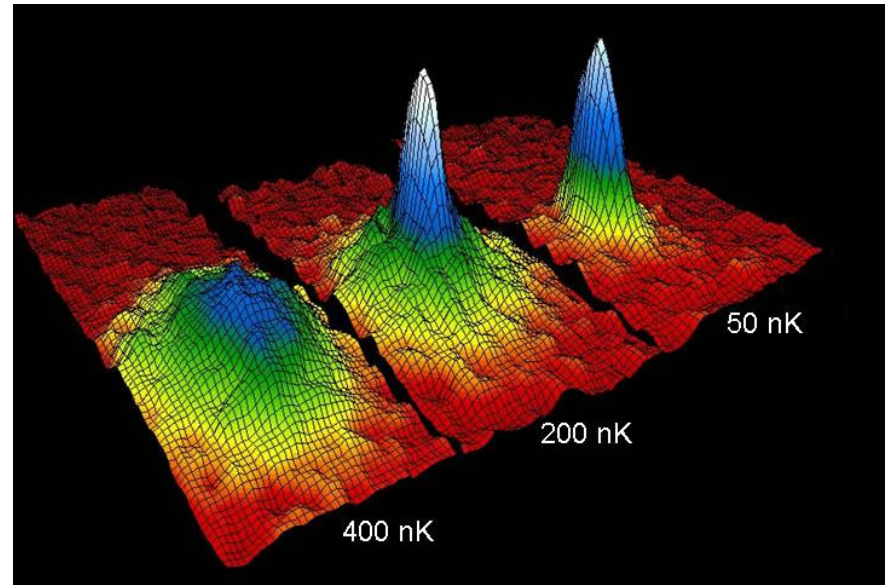


Overview: Macroscopic Quantum Tunneling and Entangled States in Bose-Einstein Condensates

- Introduction
- BEC in a double well
 - ✚ Two mode approximation
 - Potential decoherence
 - Tunneling resonances
 - ✚ Four mode approximation
 - Role of higher levels in each well
- Macroscopic Quantum Tunneling
 - ✚ Quantum sloshing in a tilted double well
- Conclusions and Outlook

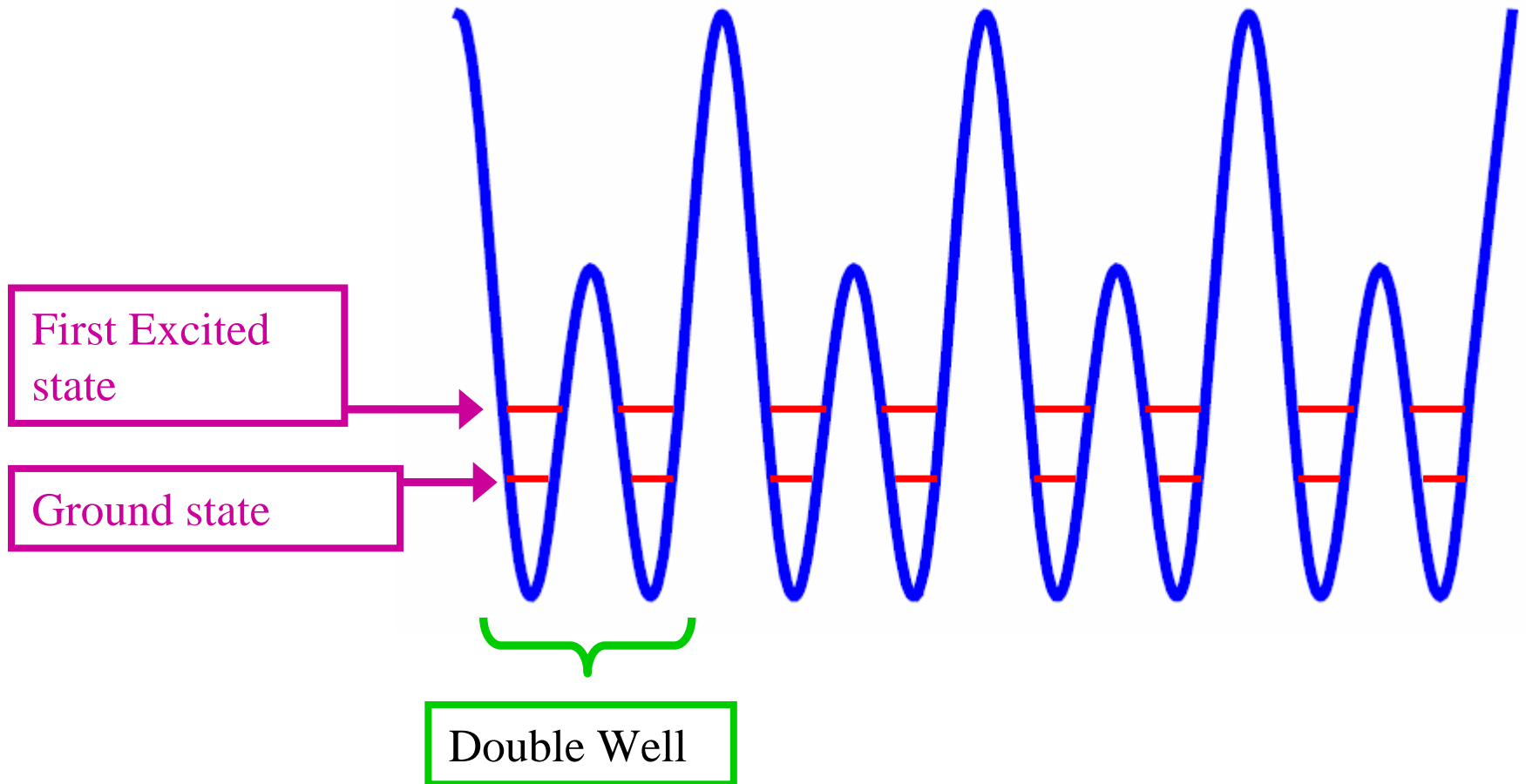
Bose-Einstein Condensation in Optical Lattices

- Combine Bose-Einstein condensate (BEC)
 - ✚ JILA: 10^6 gaseous ^{87}Rb atoms in a harmonic trap
 - Nobel prize, 2001
- With a light crystal formed by laser standing waves in 2D/3D



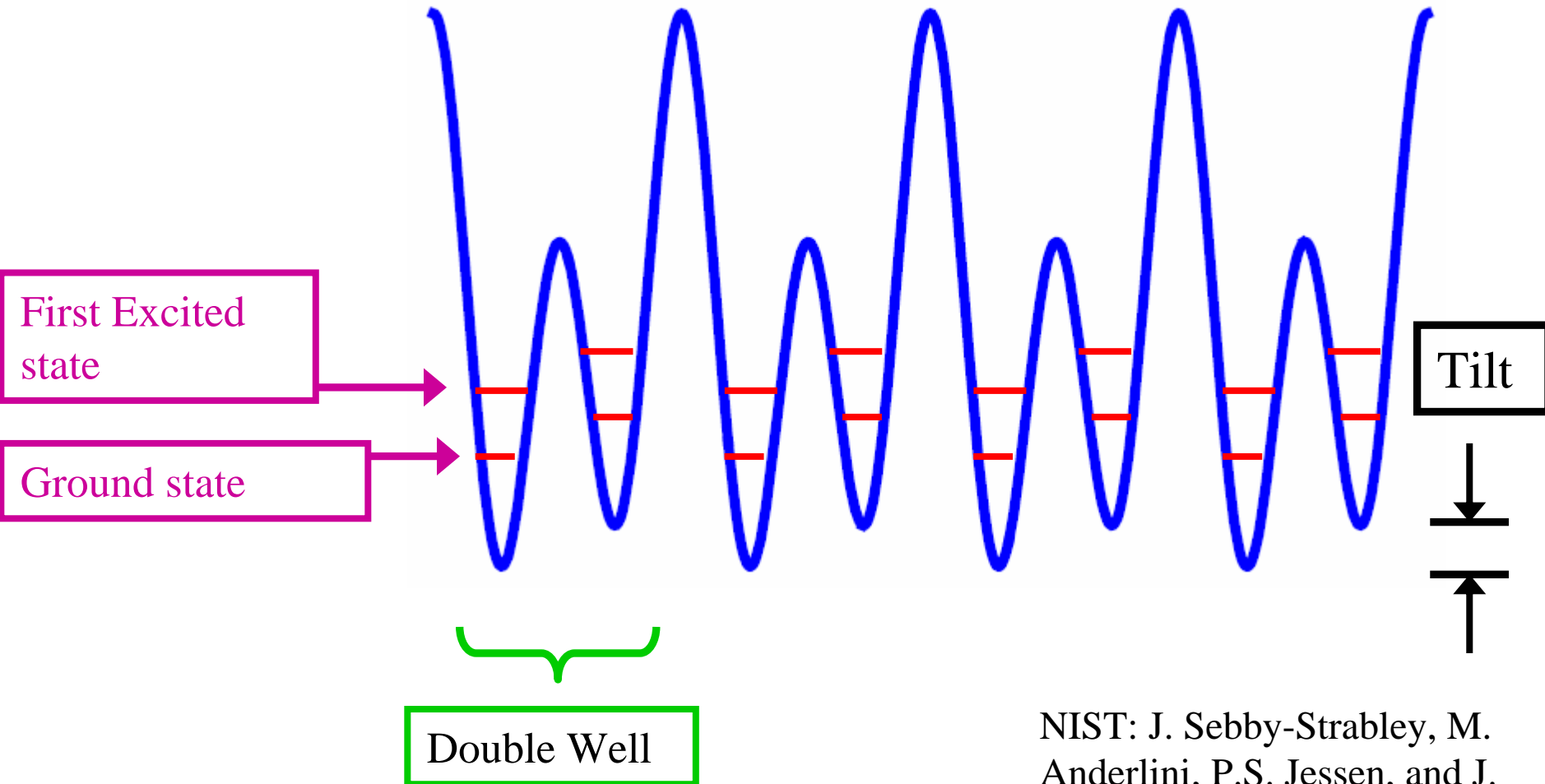
Lattice of Double Wells

- With two frequencies, we obtain



Lattice of *Tilted* Double Wells

- With more experimental cleverness, one obtains



NIST: J. Sebby-Strabley, M. Anderlini, P.S. Jessen, and J. V. Porto, PRA in press (2006)

Motivation

- Quantum many body theory of the tilted double well

- ✚ Controlled 2-qubit gates for quantum computing

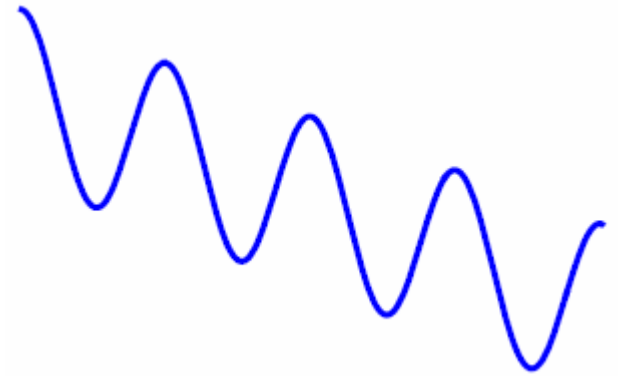
- ✚ Gravitometry, atom laser

- ✚ BEC in a double well

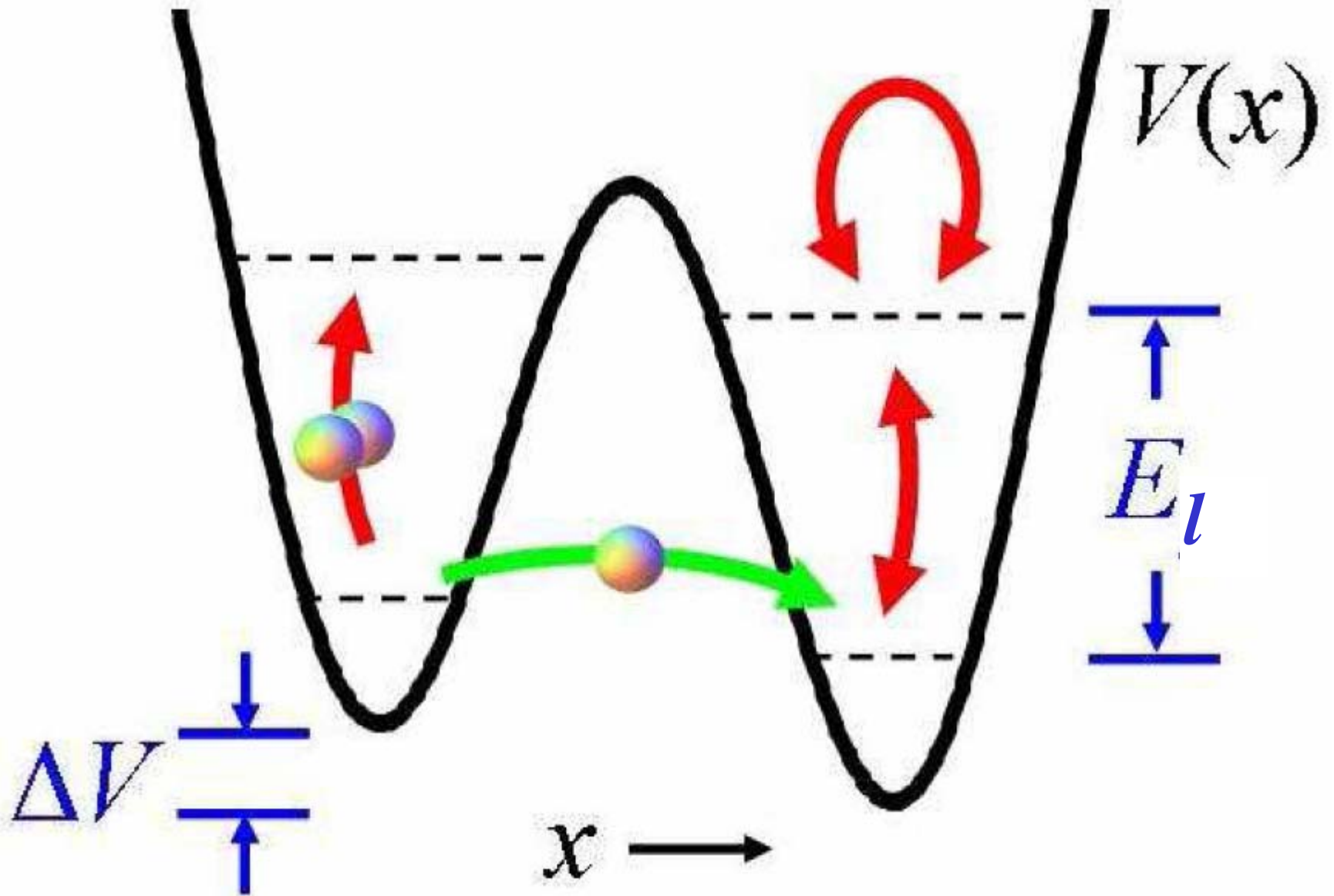
- Many-body entangled states

- ✚ Macroscopic quantum tunneling

- Push the limits of quantum mechanics



Sketch of Hamiltonian



Fundamental Equations I: Two mode Hamiltonian

- 1. Second quantized continuum field theory for weakly interacting gas in s-wave limit
- 2. Construct localized states and discretize
- 3. Obtain Bose-Hubbard-like Hamiltonian

$$\hat{H}^m = -J^m \sum_j \hat{b}_j^{m\dagger} \hat{b}_j^m + U^{mm} \sum_j \hat{n}_j^m (\hat{n}_j^m - 1) + \Delta V \hat{n}_L^m,$$

The diagram includes the following callout boxes and arrows:

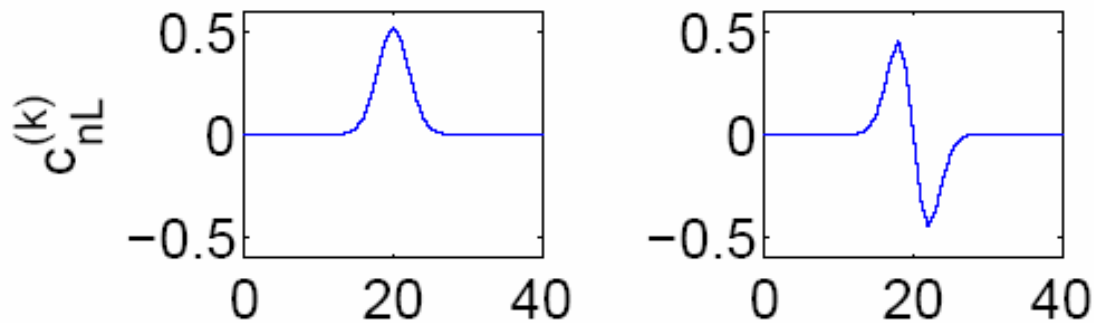
- Hopping**: Points to the $-J^m$ term.
- Level index: 0,1**: Points to the $\hat{b}_j^{m\dagger}$ and \hat{b}_j^m operators.
- Site index: L,R**: Points to the j index in the summation.
- Interaction**: Points to the U^{mm} term.
- Tilt**: Points to the ΔV term.

See, e.g., M. Jaaskelainen and P. Meystre, Phys. Rev. A 71, 043603 (2005)

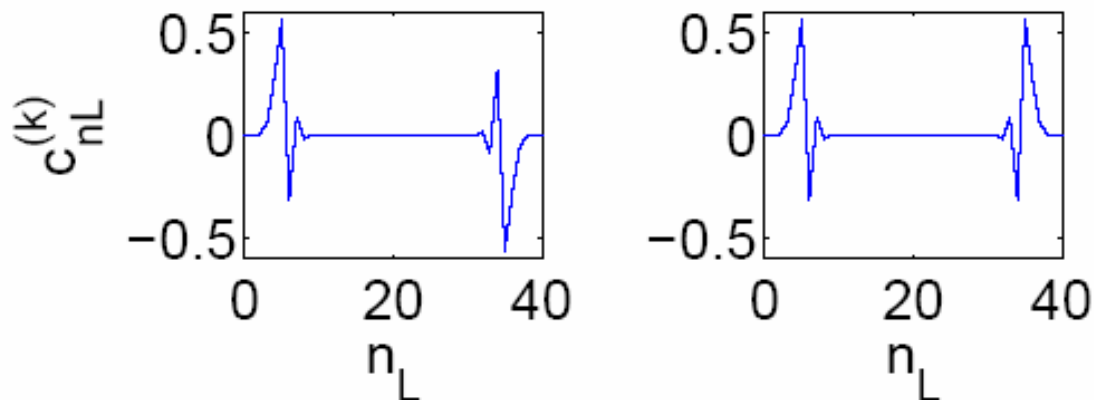
Typical solutions

- 2 sites, 1 band, no external potential, Fock states

$$|\psi^{(k)}\rangle = \sum_{n_L} c_{n_L}^{(k)} |n_L, N - n_L\rangle$$



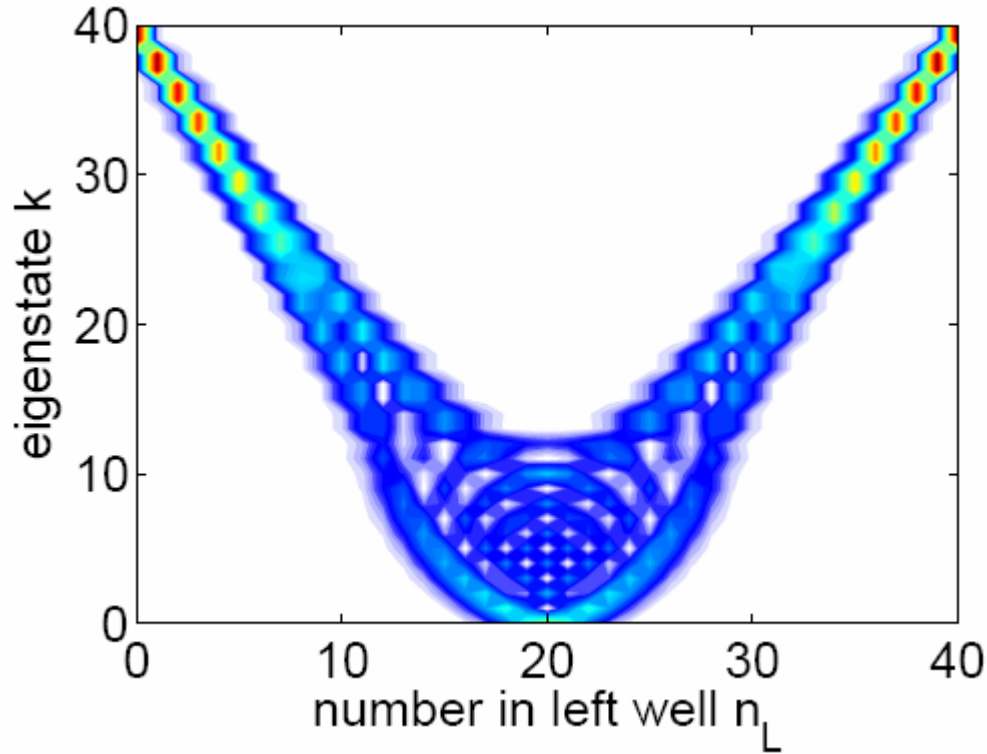
(1) Harmonic
Oscillator-like:
Superfluid to Mott



(2) Entangled:
Antisymmetric
and Symmetric

All eigenstates of two-mode Hamiltonian

- Probability density

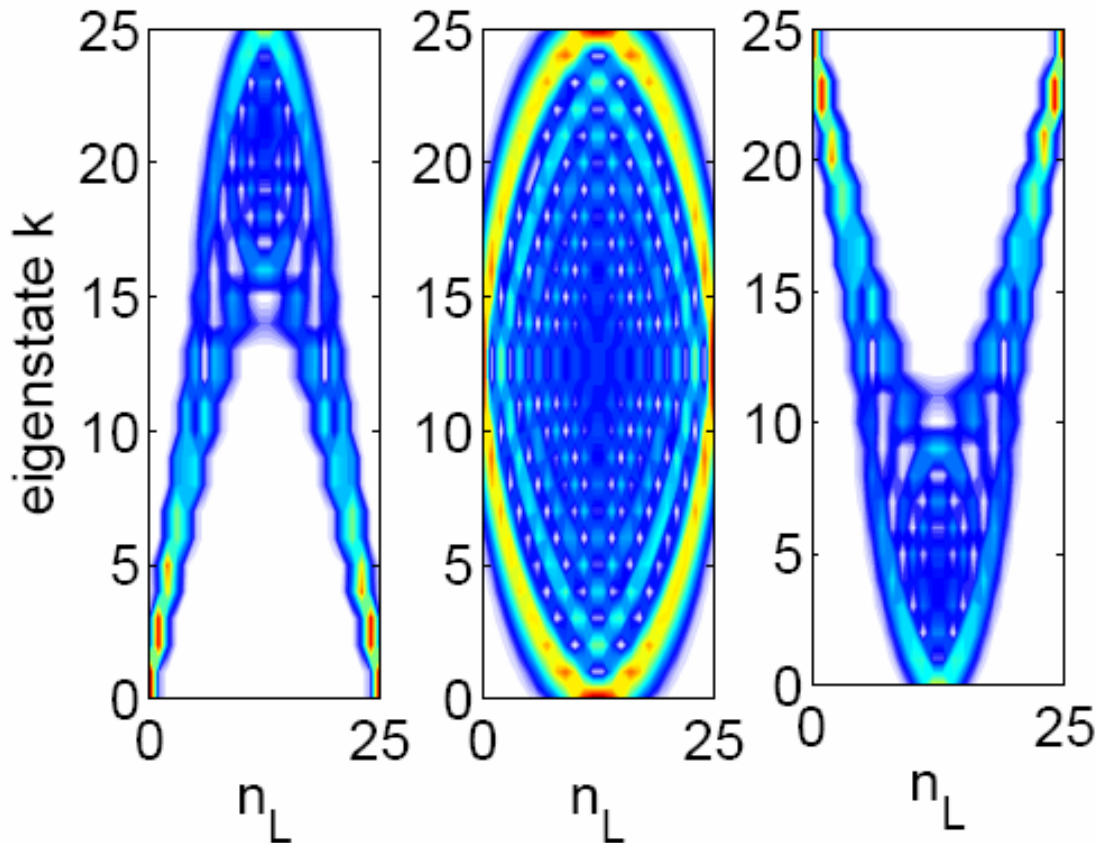


Entangled

Harmonic oscillator-like

Ratio of Hopping/Interaction

- Probability Density



Strong Attractive
Interactions

Weak
Interactions

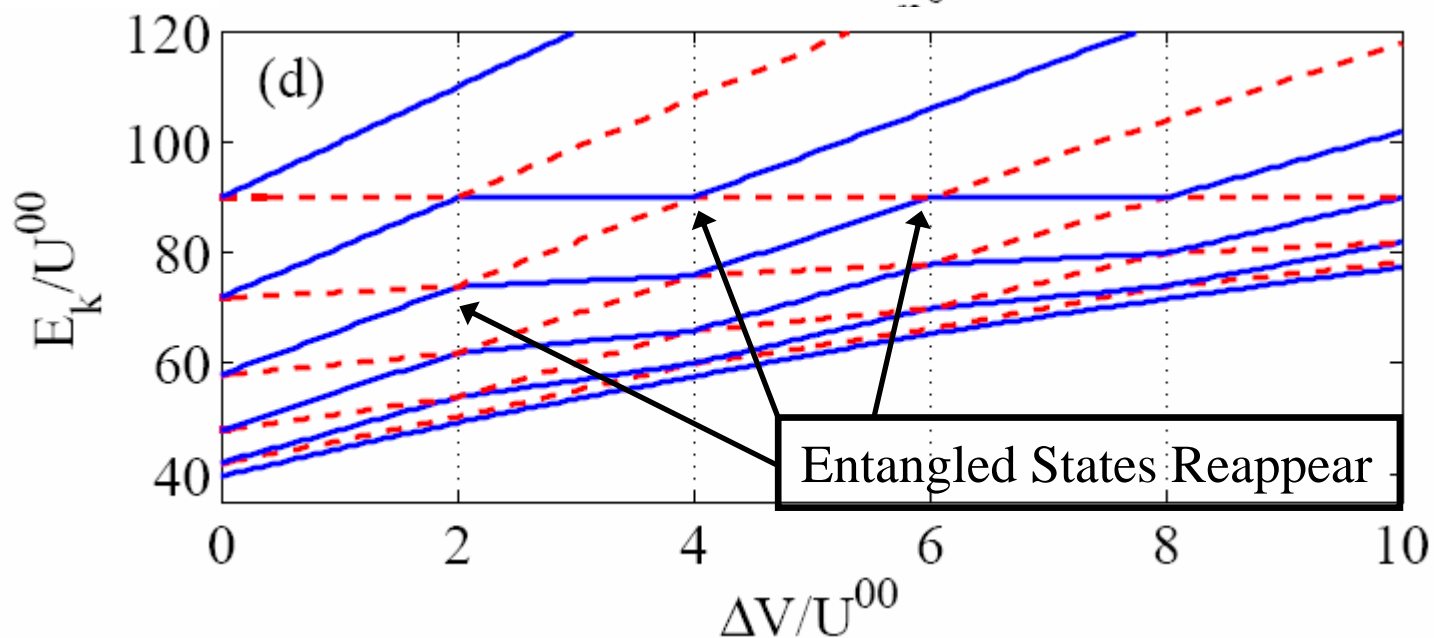
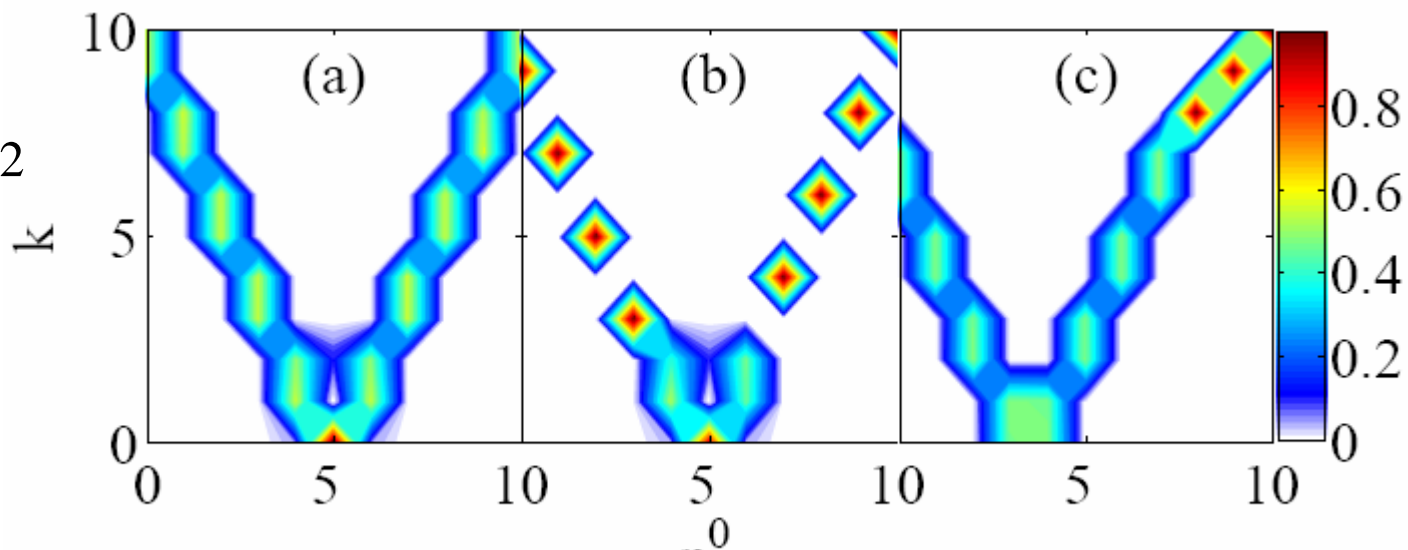
Strong Repulsive
Interactions

Response to Tilt: $N=10$

(a) No tilt

(b) $\Delta V/U=10^{-2}$

(c) $\Delta V/U=6$



Tunneling resonances

- Entangled states are fragile with respect to *potential decoherence*

- Reappear when

$$\Delta V = \Delta V_n^m \equiv 2nU^{mm}, \quad n = 1, 2, \dots, N^m - 1$$

- Width of reappearance

= width of avoided crossing

= energy difference ΔE between symmetric/anti-symmetric pairs of states

$$\Delta V_{\text{crit}}^m \equiv \frac{2\Delta E_{n_L}^m}{N^m - 2n_L^m}, \quad 0 \leq n_L^m < N^m/2$$

Fundamental Equations II: Four mode Hamiltonian

- Two-mode Hamiltonian for ground and first excited levels plus coupling terms

$$\hat{H} = \sum_m \hat{H}^m + E_\ell \hat{N}^1 \leftarrow \text{Level energy difference}$$

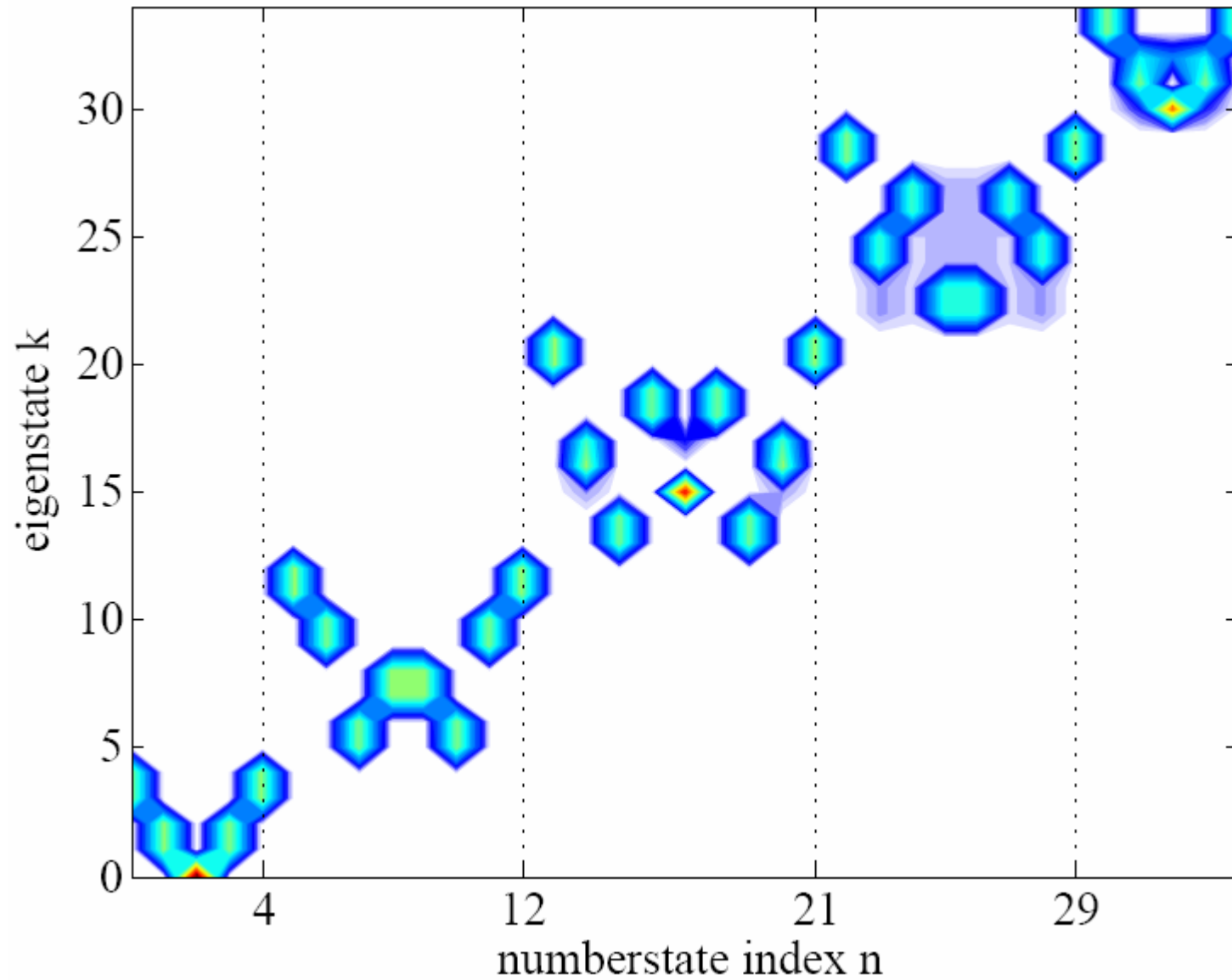
$$+ U^{01} \sum_{j, m \neq m'} \left[2\hat{n}_j^m \hat{n}_j^{m'} + \hat{b}_j^{m\dagger} \hat{b}_j^{m\dagger} \hat{b}_j^{m'} \hat{b}_j^{m'} \right]$$

Inter-level interaction
Inter-level pair hopping

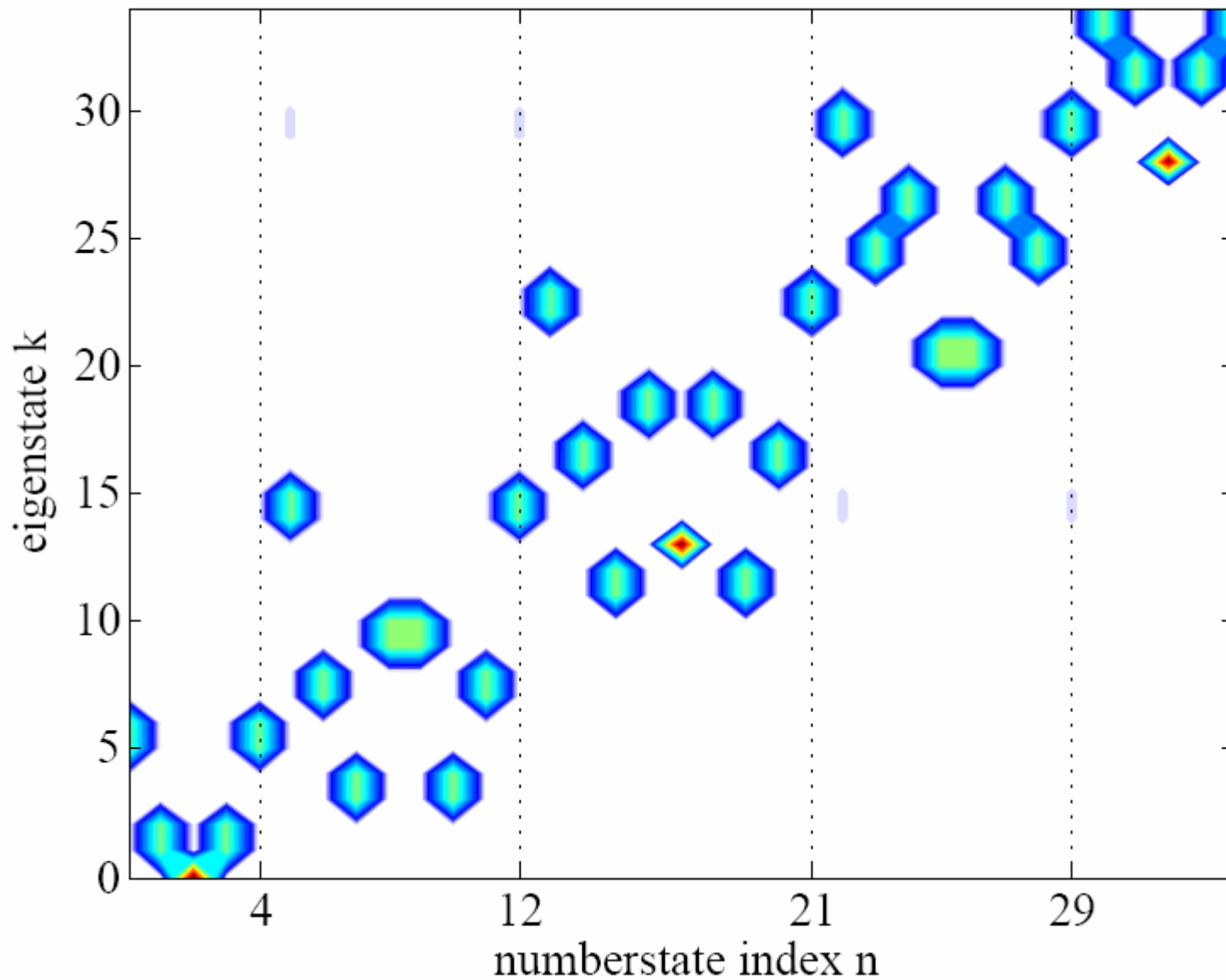
- Fock state Hilbert space

$$|\Psi\rangle = \sum_n c_n |n\rangle, \quad |n\rangle = |n_L^0, n_R^0\rangle \otimes |n_L^1, n_R^1\rangle$$

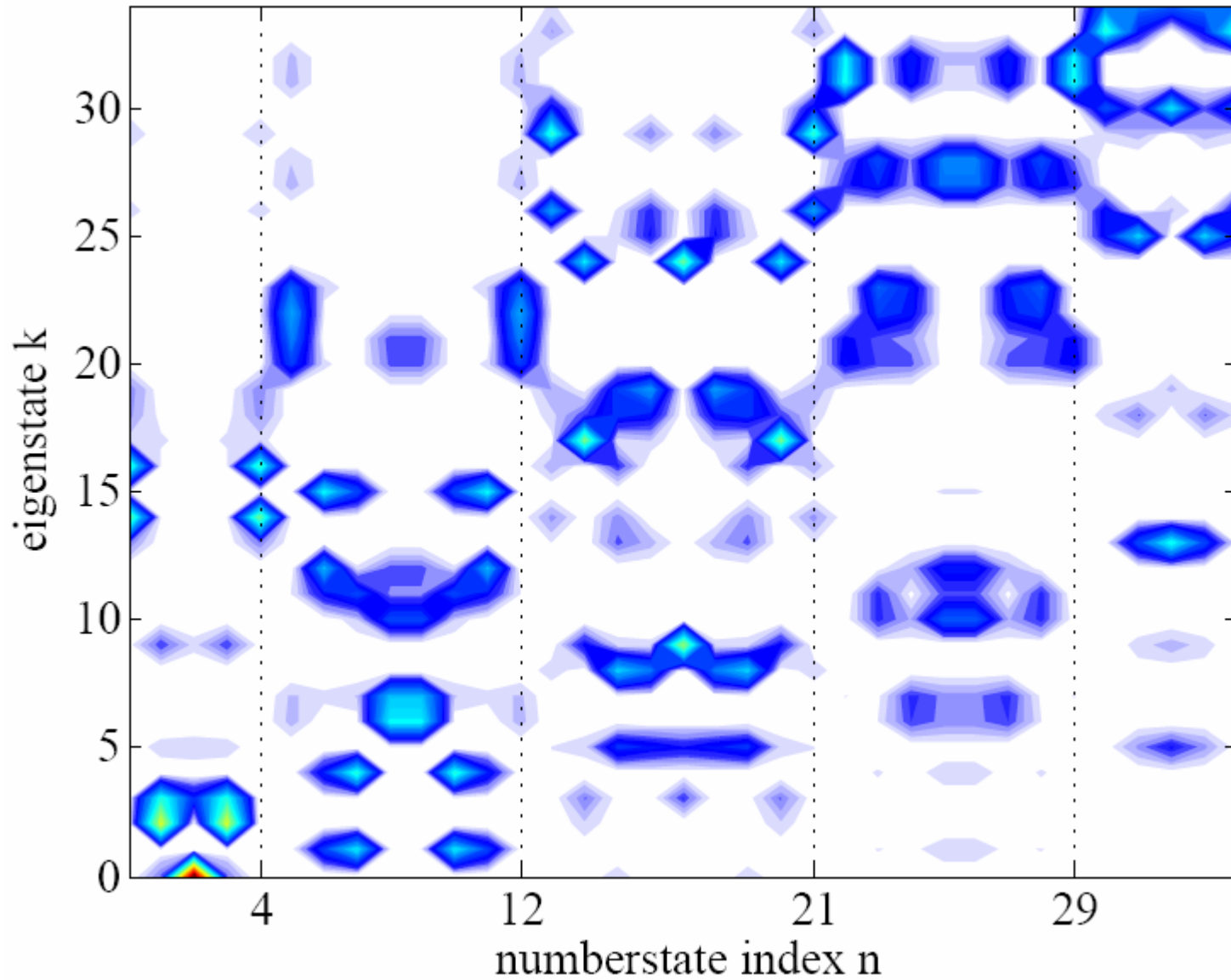
Four-mode Model Stationary States I



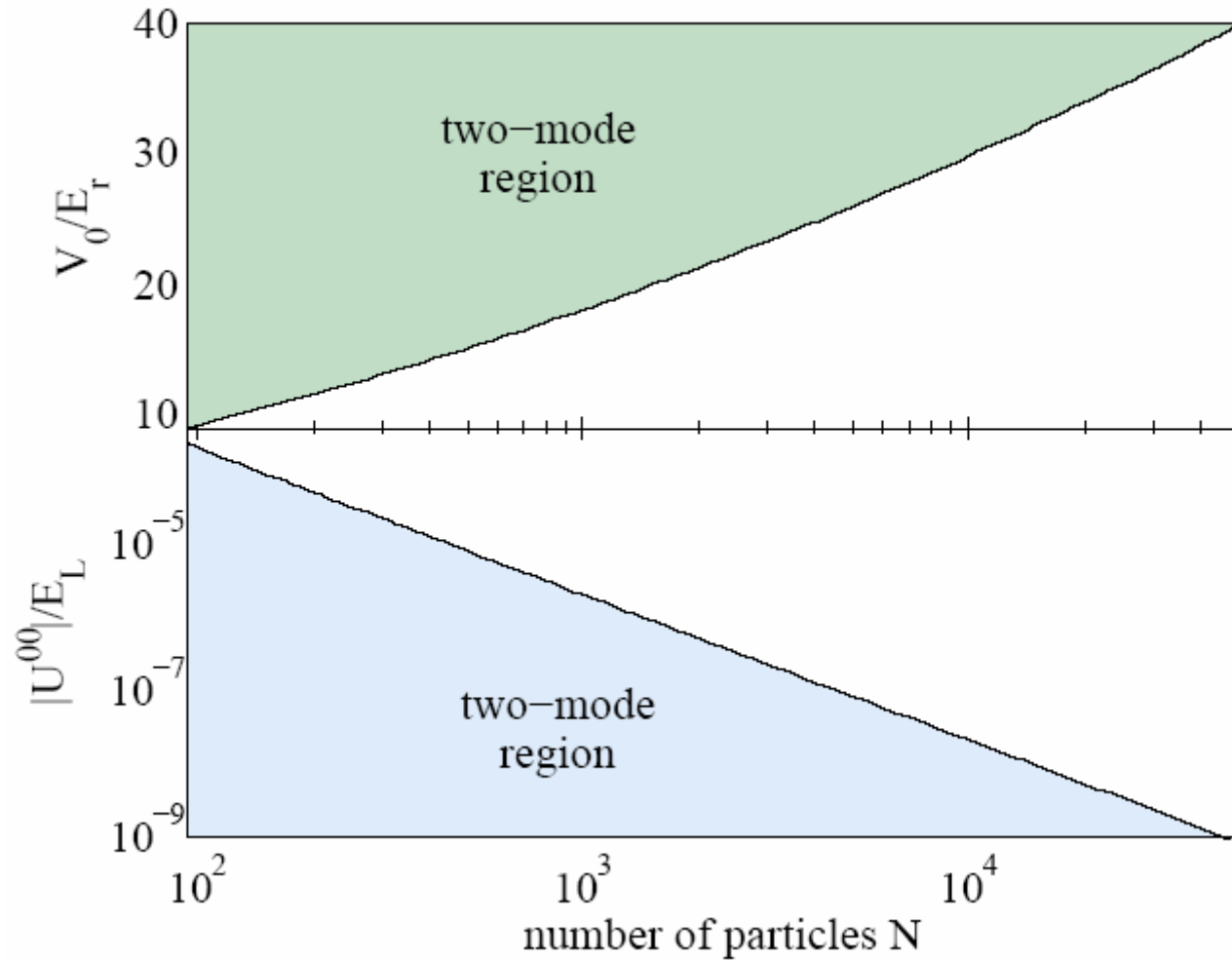
Four-mode Model Stationary States II



Four-mode Model Stationary States III



Bounds on Use of Two-mode Approximation

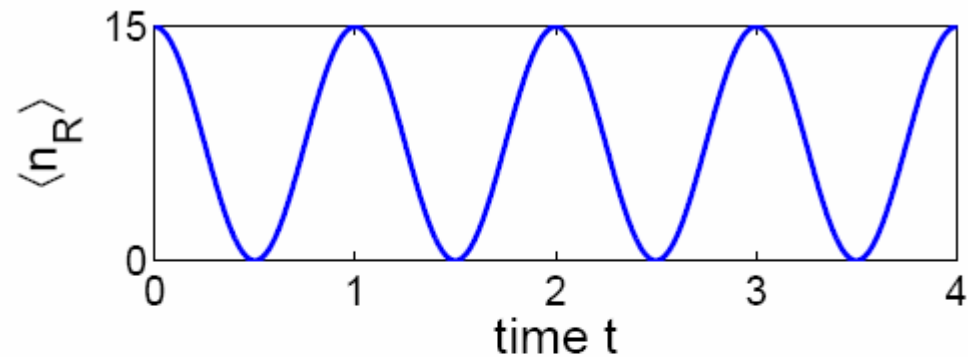


Sloshing in a Tilted Double Well: The Basic Idea

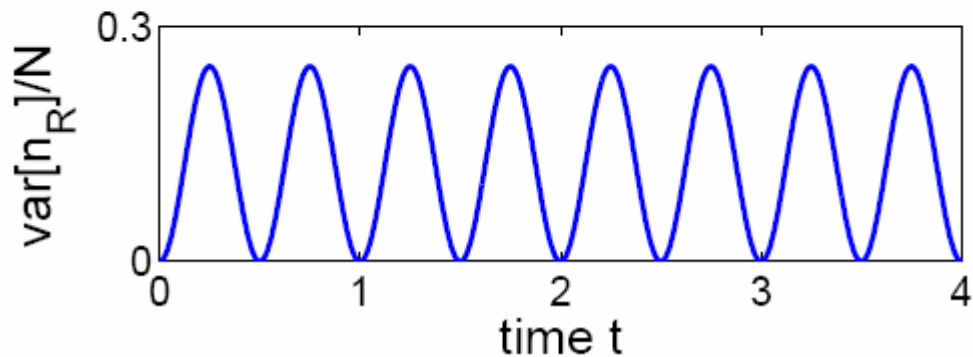
- Build up a picture of MQT particle by particle
- Tune towards Mott border (raise lattice barrier height)
- Observable in experiments

Dynamics: Non-interacting ($U=0$)

- Average number and average normalized variance

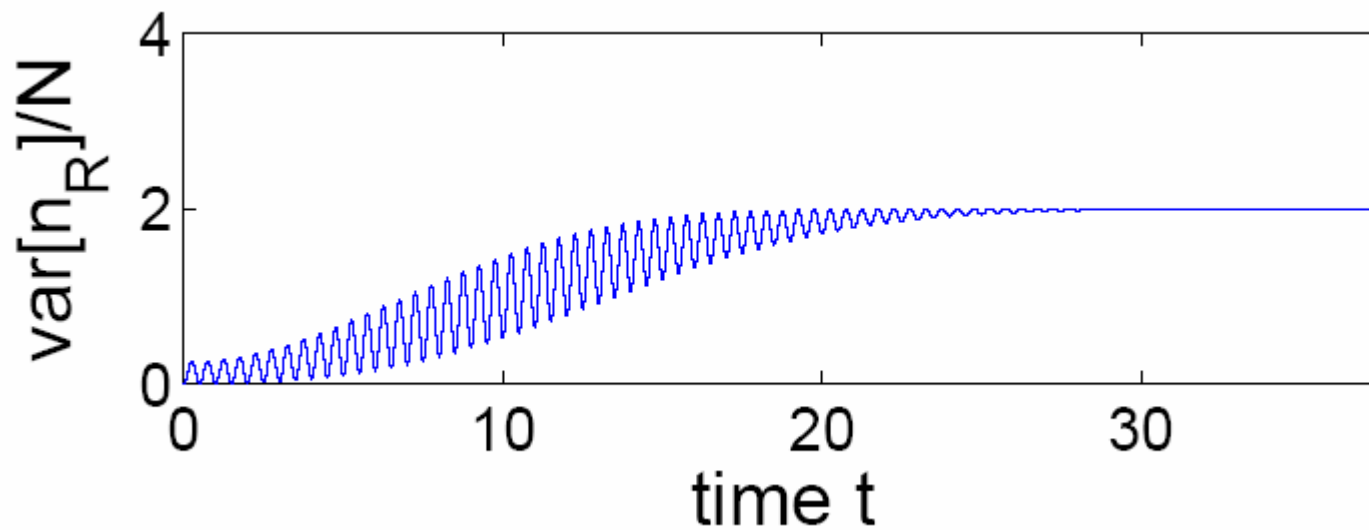
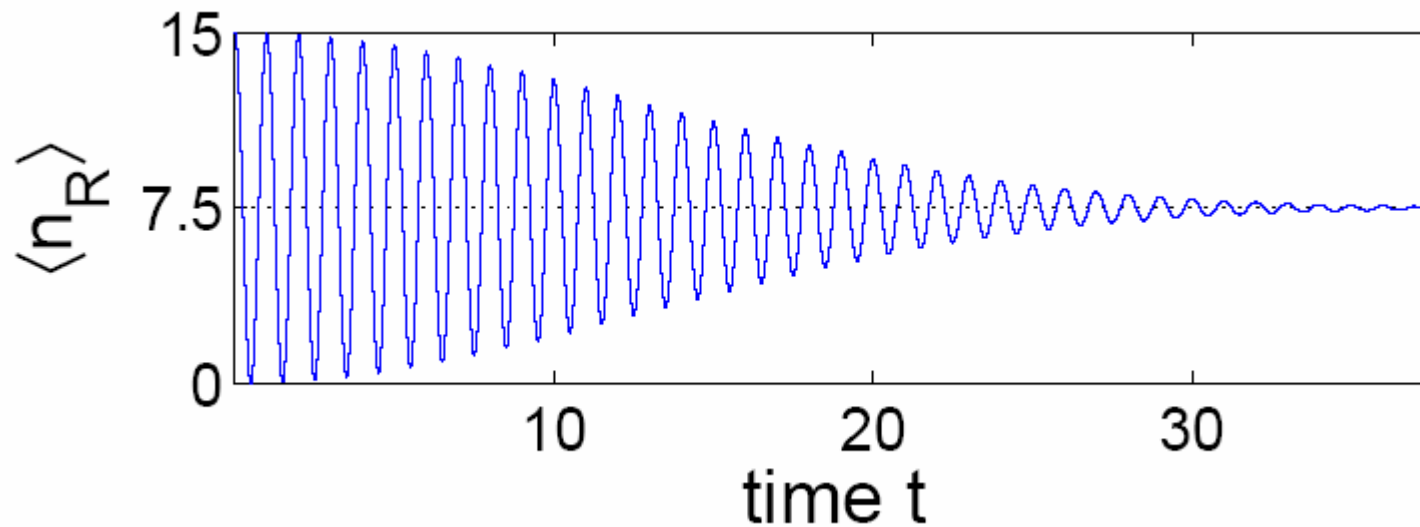


Time in units of
 $\pi\hbar/J$



$$\langle \hat{n}_R \rangle = N(1 + \cos 2Jt/\hbar)$$

Weak Interactions ($J/NU \gg 1$), Short Time Oscillations

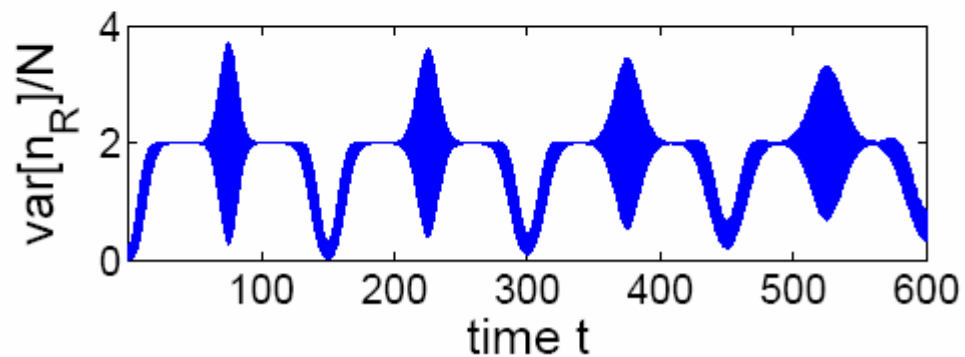
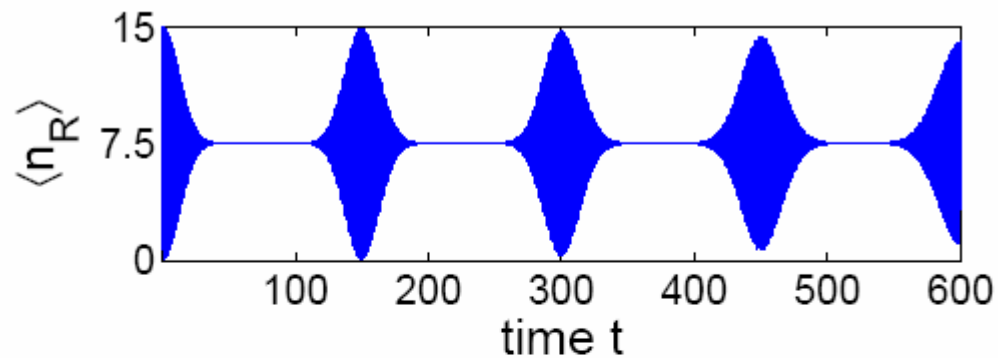


Weak Interactions: Long Time Oscillations

- $\langle \hat{n}_R \rangle = N \left(1 + \cos(2Jt/\hbar) \cos(Ut/\hbar)^{N-1} \right)$

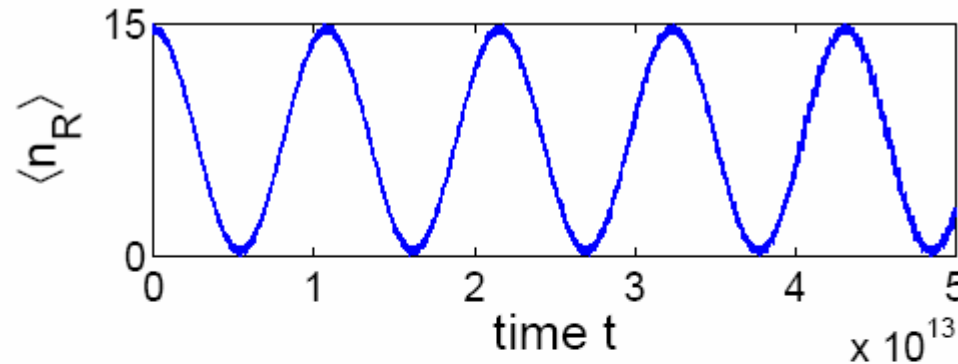
⊕ Tunneling \rightarrow carrier frequency

⊕ Interactions \rightarrow envelope frequency

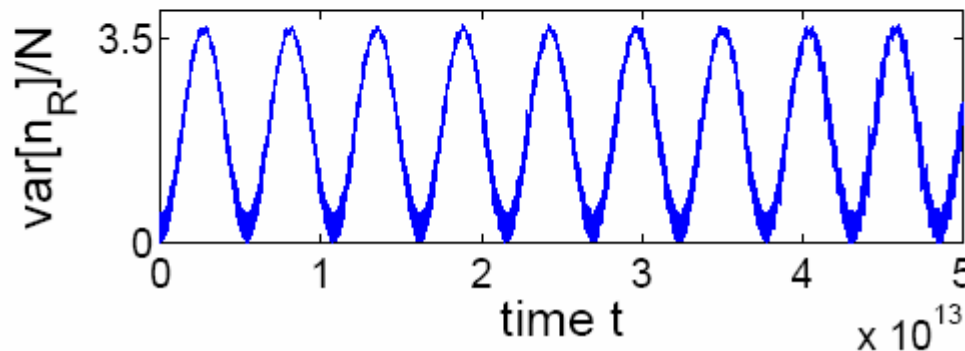


Strong Interactions/High Barrier ($J/NU \ll 1$)

- $\langle \hat{n}_R \rangle = N (1 + \cos(\Delta E_N t / \hbar))$
 - ✚ ΔE_N is the splitting between antisymmetric and symmetric extreme cat states $|0, N\rangle \pm |N, 0\rangle$



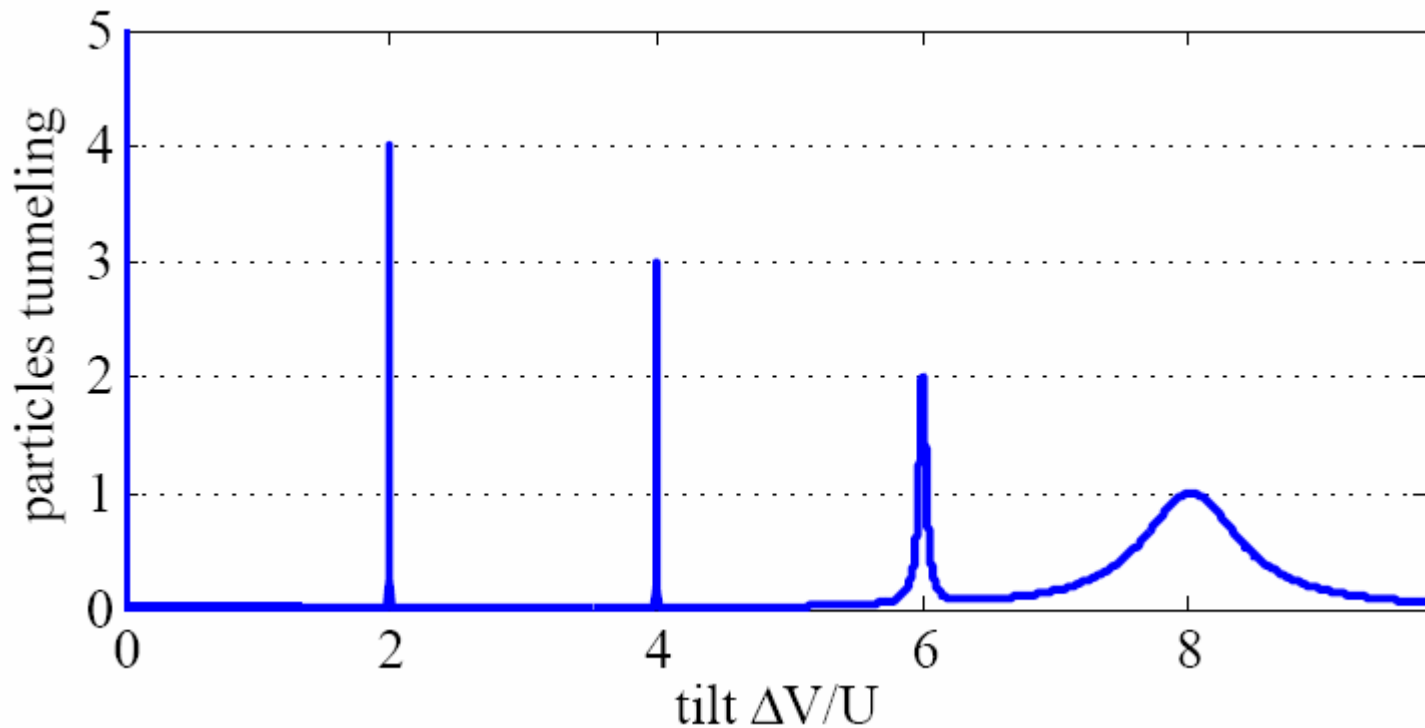
Time in units of
 \hbar/NU



$$T_{osc} \propto (U/J)^{N-1} \pi \hbar / J$$

Tunneling Resonances

- Tilt $2Uk \rightarrow$ N-k particles slosh



Conclusions

- Higher mode effects cannot be neglected in many experiments
- Entangled states are fragile to potential decoherence
 - ✚ Reappear for tilt proportional to interaction strength
- MQT tunneling times exponential in number of particles in high barrier limit