

Tunable negative refraction based on quantum interference

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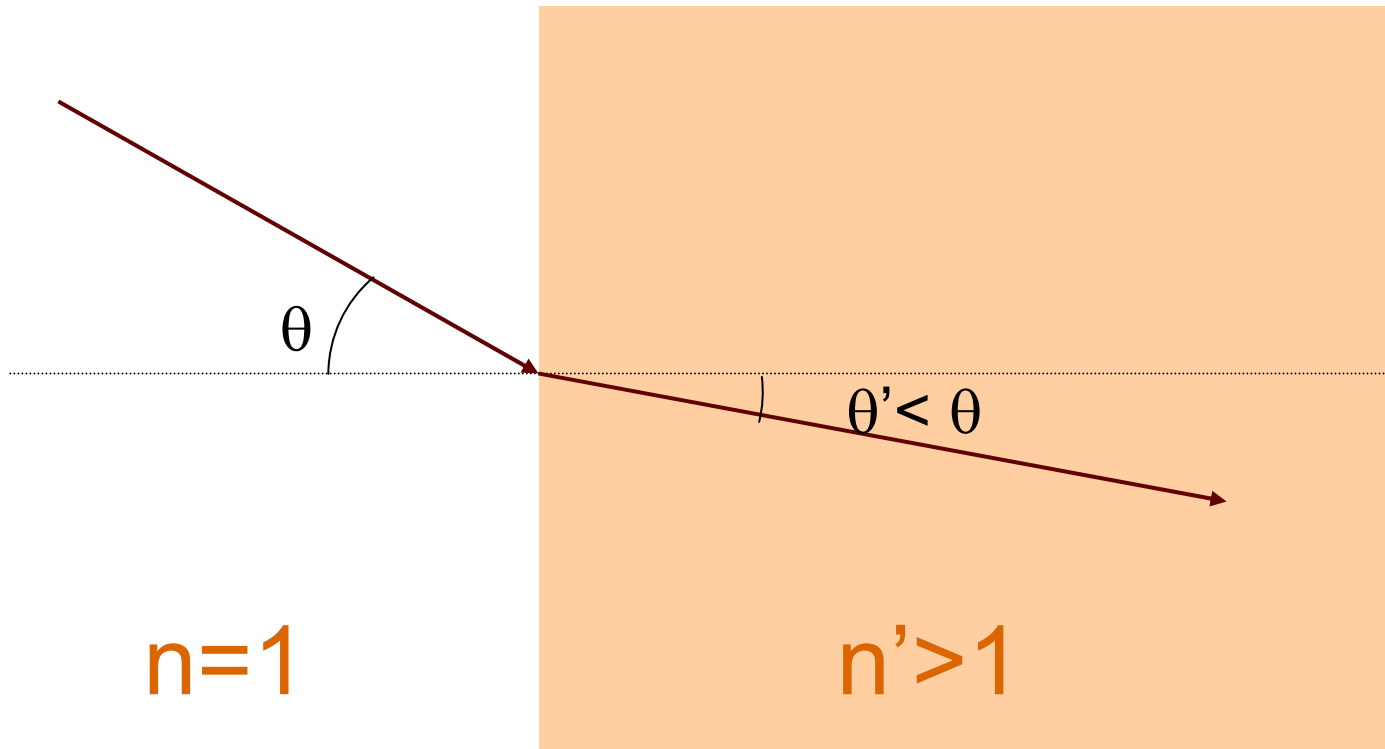
TU Kaiserslautern, Germany

Ron Walsworth

Harvard/Smithsonian CfA

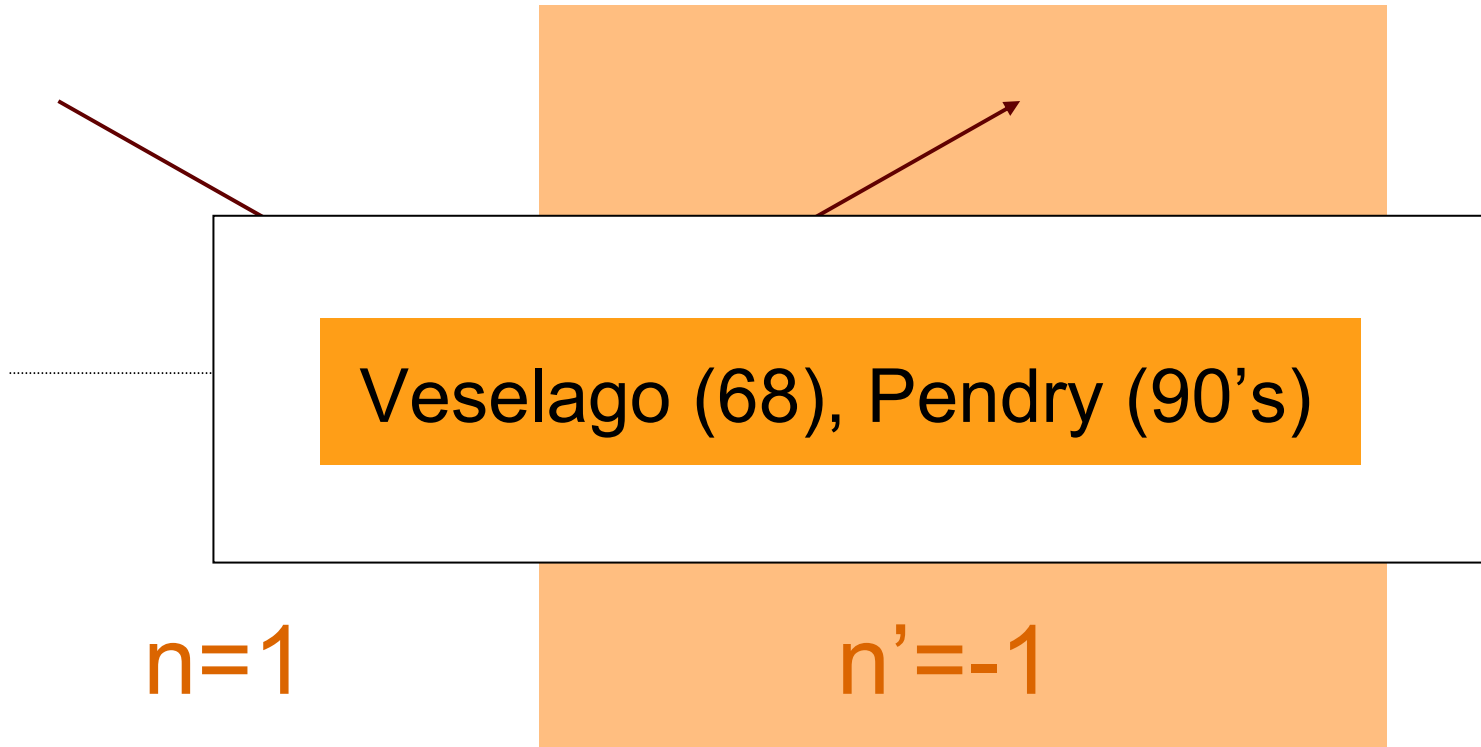
Breckenridge, August 25, 2006

Model



normal refraction
(Snell's law)

Model

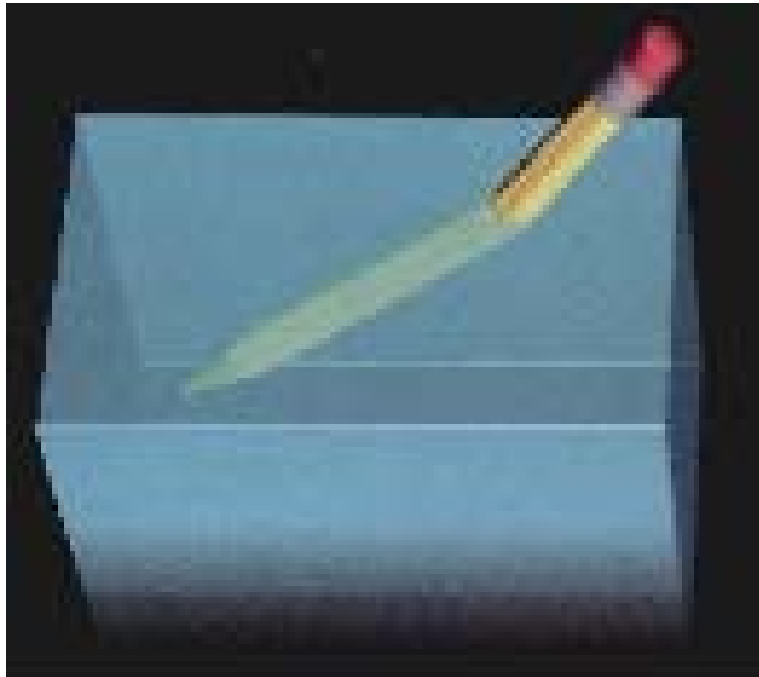


negative refraction
(Snell's law)

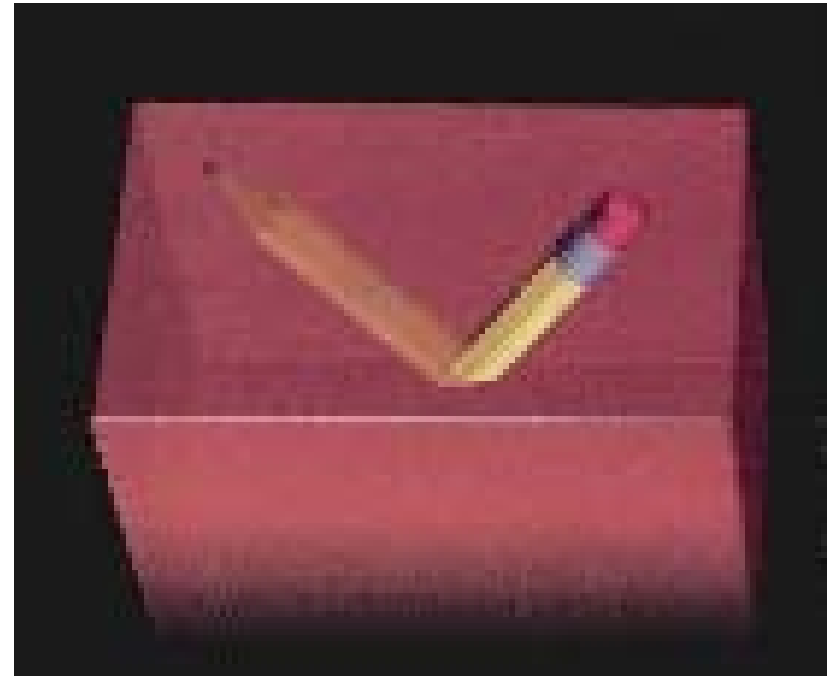


“left-handed”:
E, **B**, and **k**

Scientific American, July 2006



normal refraction



negative refraction

Contents

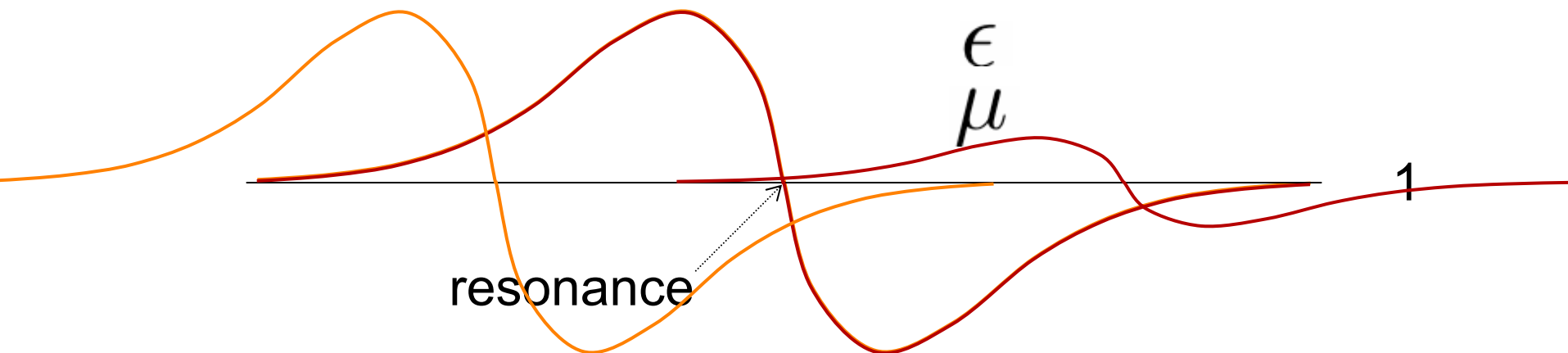
- Definitions & examples
- Chiral media
- **EIT based negative refraction**
 - Local field effects
 - Tunability
- Conclusion

Origin of negative refraction

$$n = \sqrt{\epsilon\mu}$$

ϵ : electric permittivity
 μ : magnetic permeability

With both, ϵ and μ negative n negative



Applications and Definitions

- Applications:
 - perfect lens



Superlens

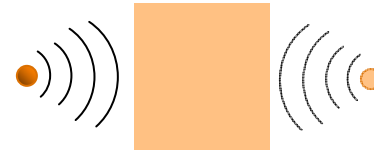
normal lens:

$$k_x = \sqrt{\frac{\omega^2}{c^2} - k_z^2}$$

☑ Resolution: $2\pi k_x^{-1} \geq \lambda !$

Applications and Definitions

- Applications:
 - perfect lens



- Definitions:

- negative refraction
- left-handed materials
- chiral materials
- meta-materials

$$\text{Re}(n) < 0$$

E, **B**, and **k** are lefthanded

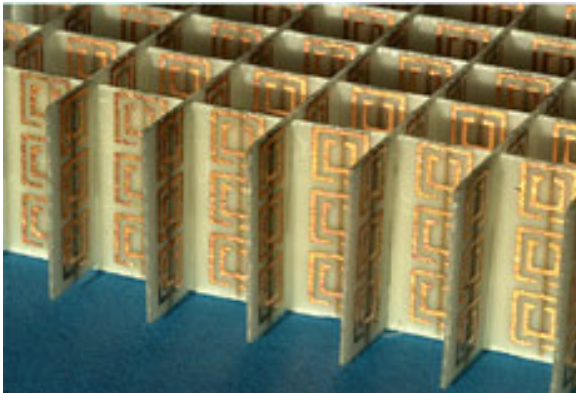
Turned polarization \Leftrightarrow **E**-

B-cross coupled

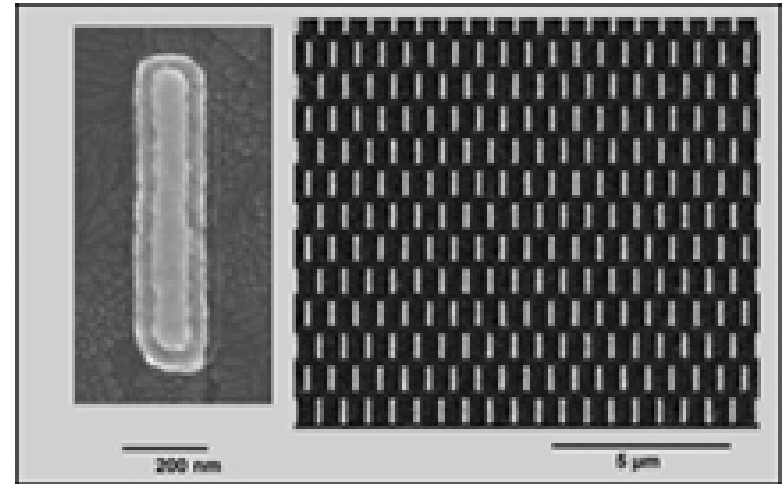
man-made refraction

Material examples

- μ -wave structures:



Pendry



Shalaev

Photonic bandgap material

- Use band structure of the photonic crystal to get a left-handed material (“flip over” k vector direction on Fermi surface)

QuickTime™ and a
TIFF (Uncompressed) decompressor
are needed to see this picture.

1 μ m

Notomi

- For certain frequency: **negative refraction**
- But: not “metamaterial”: No resolution beyond λ !
(**no superlensing!**)

Absorption

So far: refraction/absorption $\approx 1 \dots 5$

Our case: $\text{Re}(n)/\text{Im}(n) = 100$

Occurrence of negative refraction

- Why does negative index not occur in Nature?

Large χ_m very difficult to achieve!

Optical frequencies

Magnitude of χ_m :

$$|\chi_m| \approx \left(\frac{\mu_{\text{atom}}}{d_{\text{atom}}} \right)^2 |\chi_e| \approx \frac{1}{137^2} |\chi_e|$$

Chiral media (Pendry)

- Remember: $n = \sqrt{\epsilon\mu}$

Chiral media (Pendry)

- Remember: $n = \sqrt{\epsilon\mu}$
- Chiral media: cross coupling between electric and magnetic fields

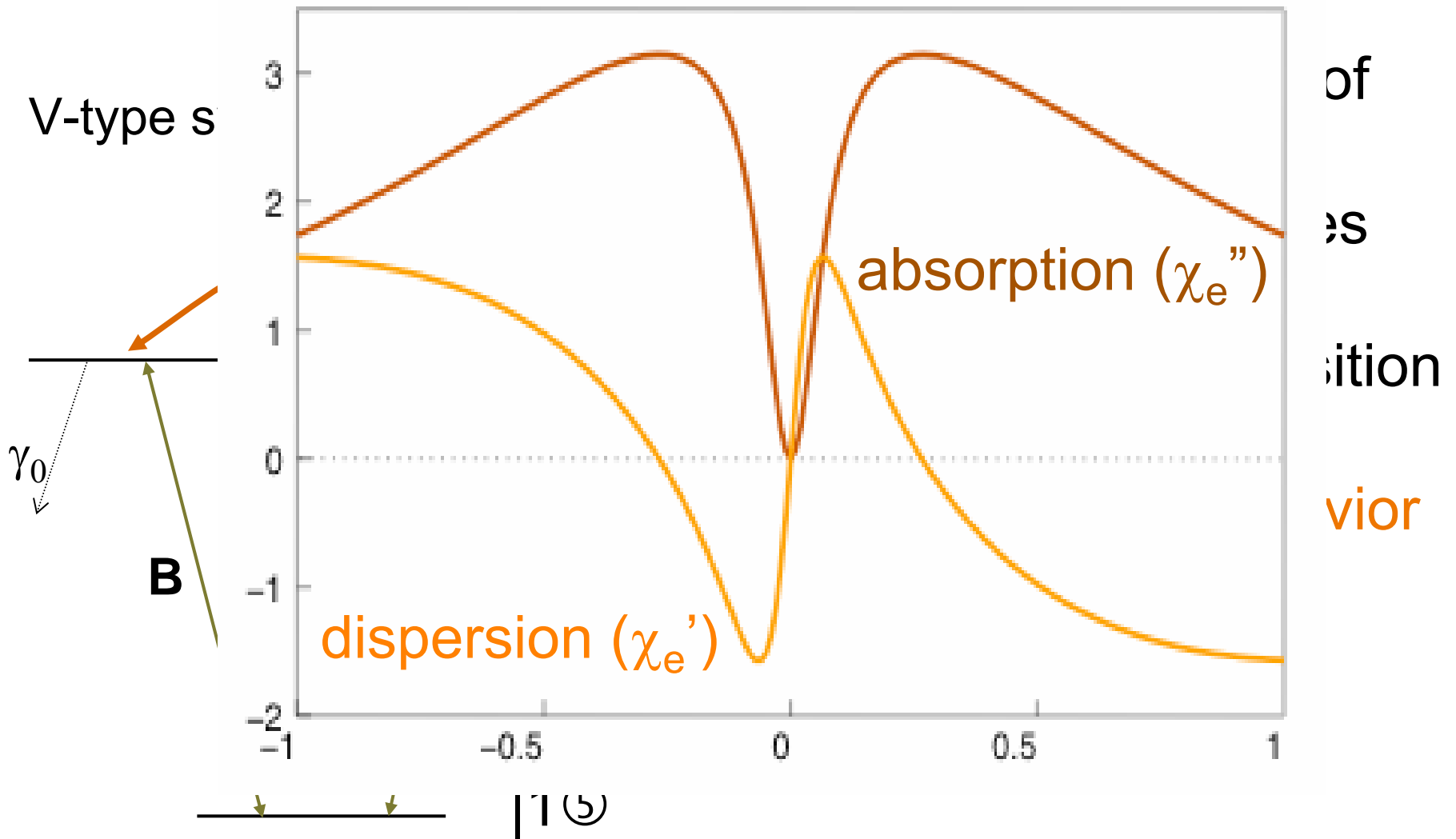
$$\begin{aligned}\mathbf{P} &= \chi_e \mathbf{E} + \xi_{eb} \mathbf{B} \\ \mathbf{M} &= \xi_{be} \mathbf{E} + \chi_m \mathbf{B}\end{aligned}\quad \text{with } |\xi| \propto \frac{1}{137} |\chi_e|$$

- Index of refraction

$$n = \sqrt{\epsilon\mu} - \xi$$

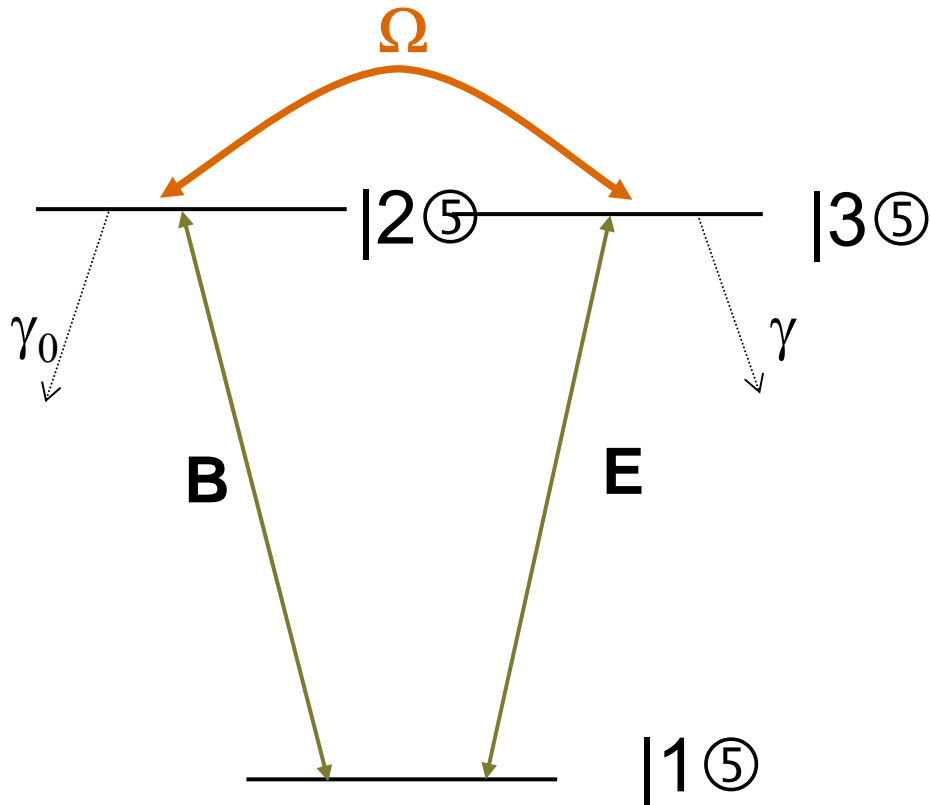
If we choose $\xi_{EH} = -\xi_{HE} = i\xi$

EIT based negative refraction



EIT based negative refraction

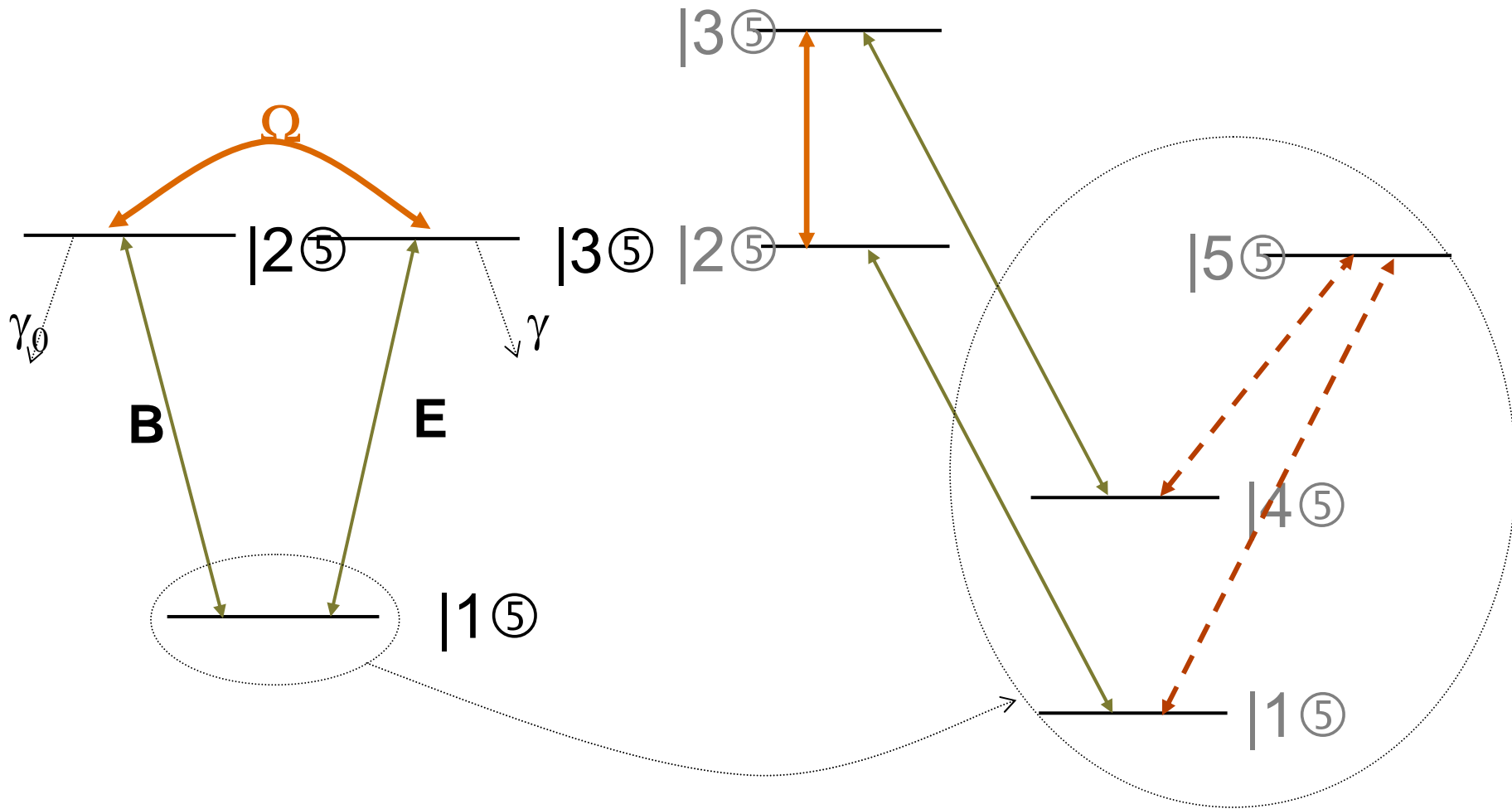
V-type system:



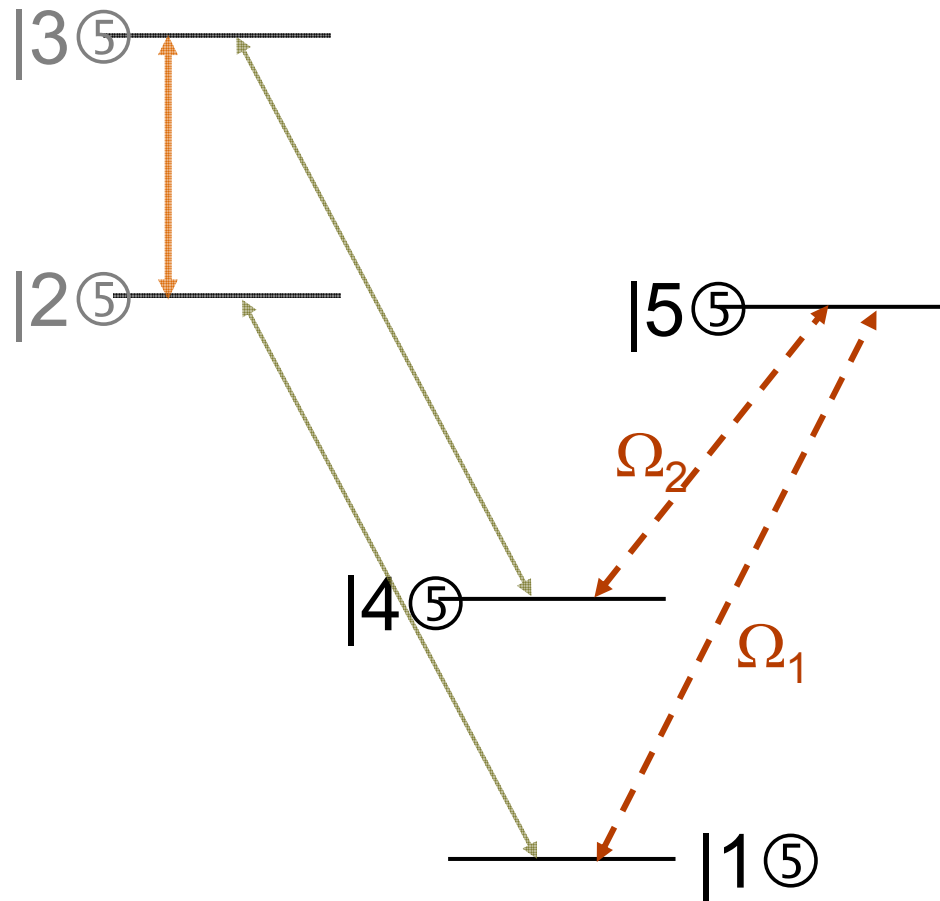
Problems:

- Ω : dc-coupling
phase of ξ not free to choose
- Ω dc-coupling: very weak Rabi frequency
- no EIT for inhomogeneously broadened systems
- level scheme hard to find in real systems

Realistic schemes



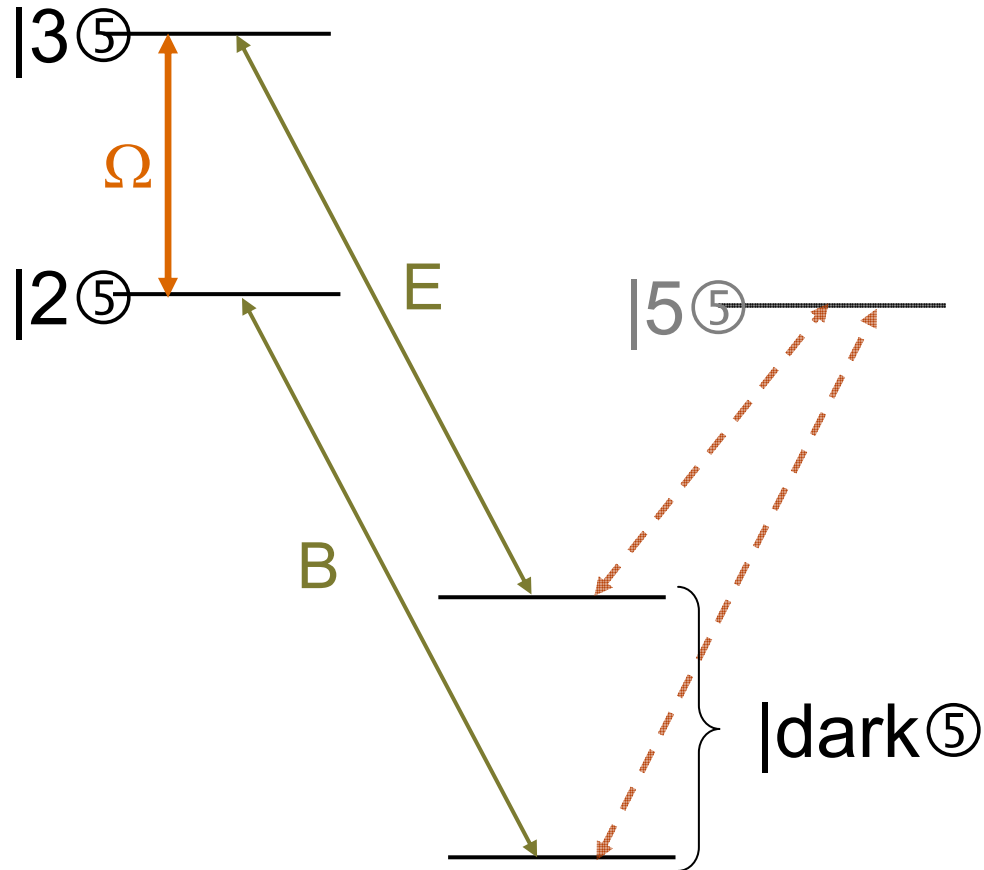
Realistic schemes



- Create dark state in superposition of $|1\rangle$ and $|4\rangle$
- Dark state acts like g.s. in 3-level system

$$|\text{dark}\rangle \propto \Omega_1 |1\rangle - \Omega_2 |4\rangle$$

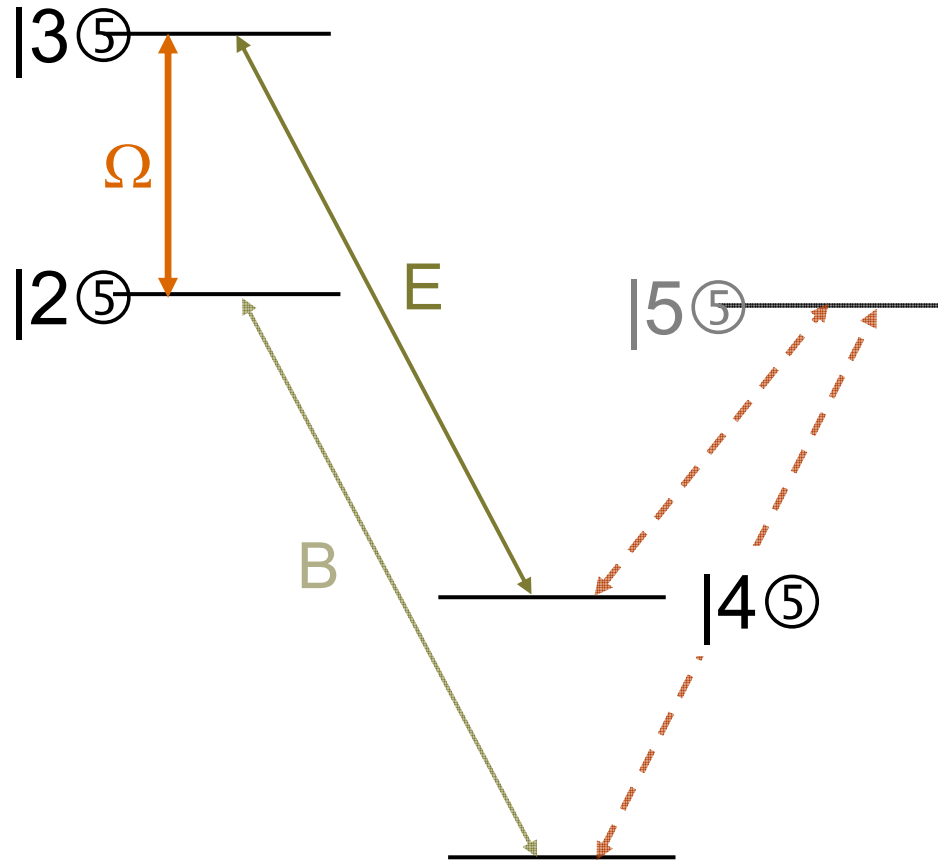
Realistic schemes



Advantages:

- Non-dc coupling field Ω
- ☑ Choose phase

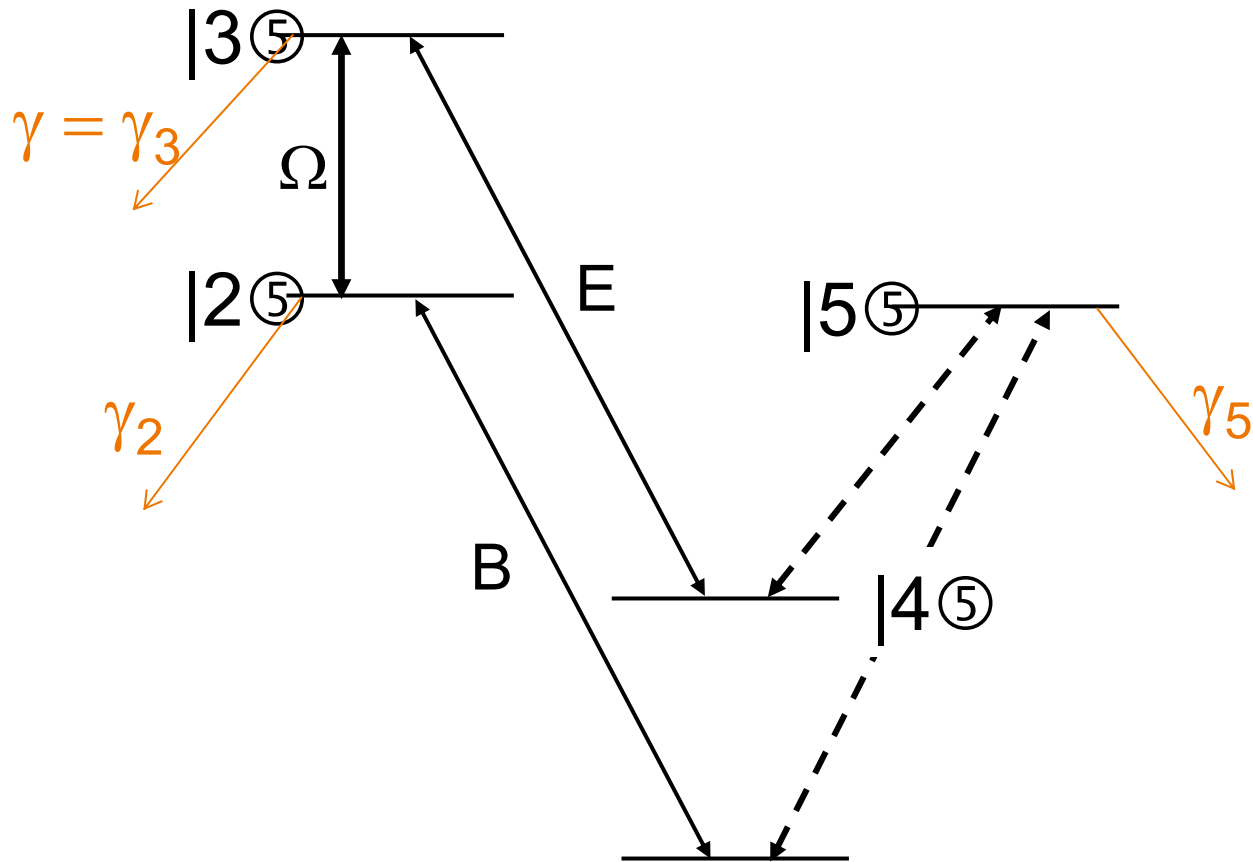
Realistic schemes



Advantages:

- Non-dc coupling field Ω
 - ☑ Choose phase
- States $|2^\oplus$ and $|4^\oplus$ can be chosen at similar energy
 - ☑ No Doppler broadening on sensitive Λ -type scheme ($|4^\oplus$, $|2^\oplus$, and $|3^\oplus$)
- Easier to realize

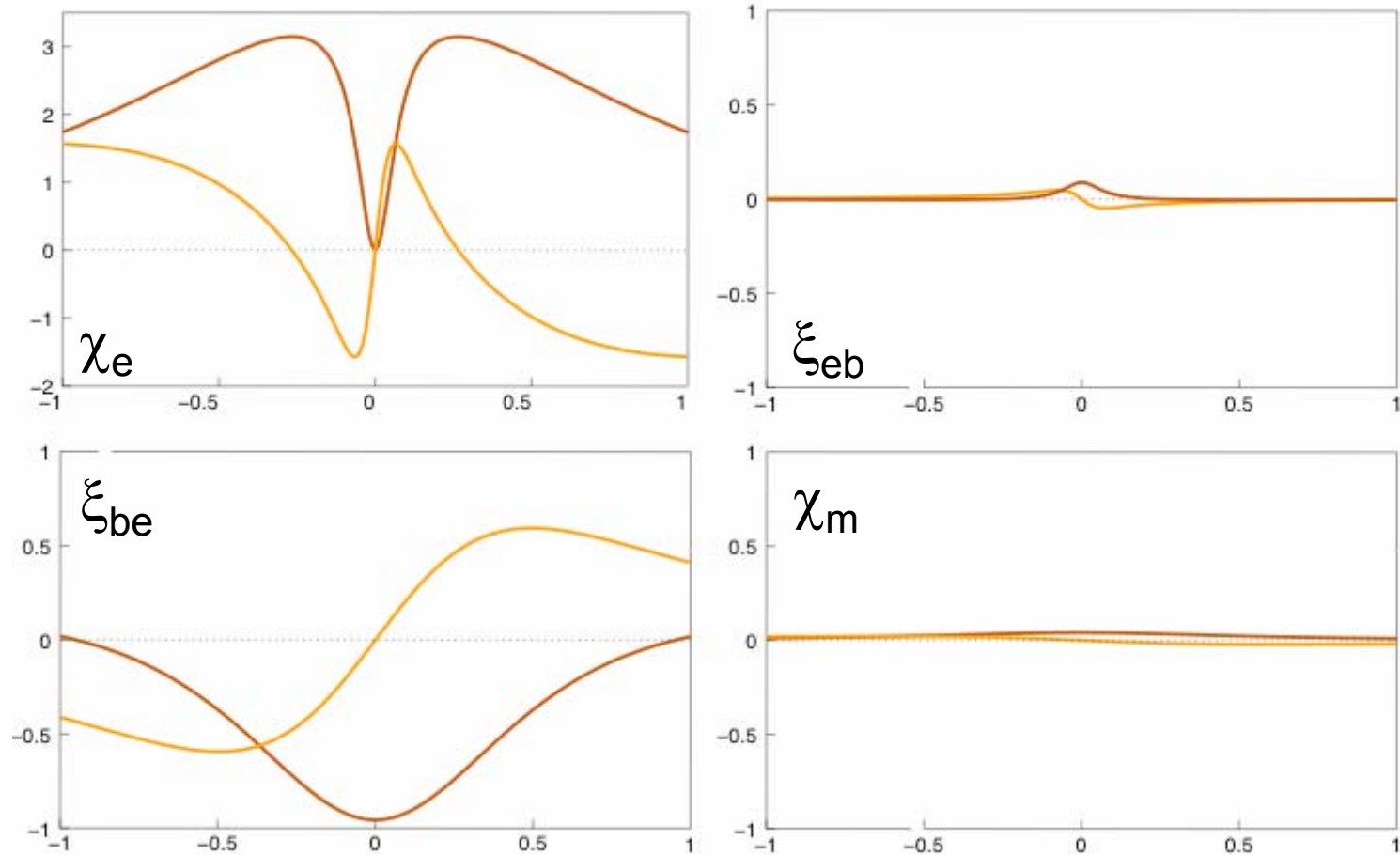
Realistic schemes



+ line broadening (inhomogeneous)

Cross couplings

Inhomogeneous broadening \approx decay rate γ

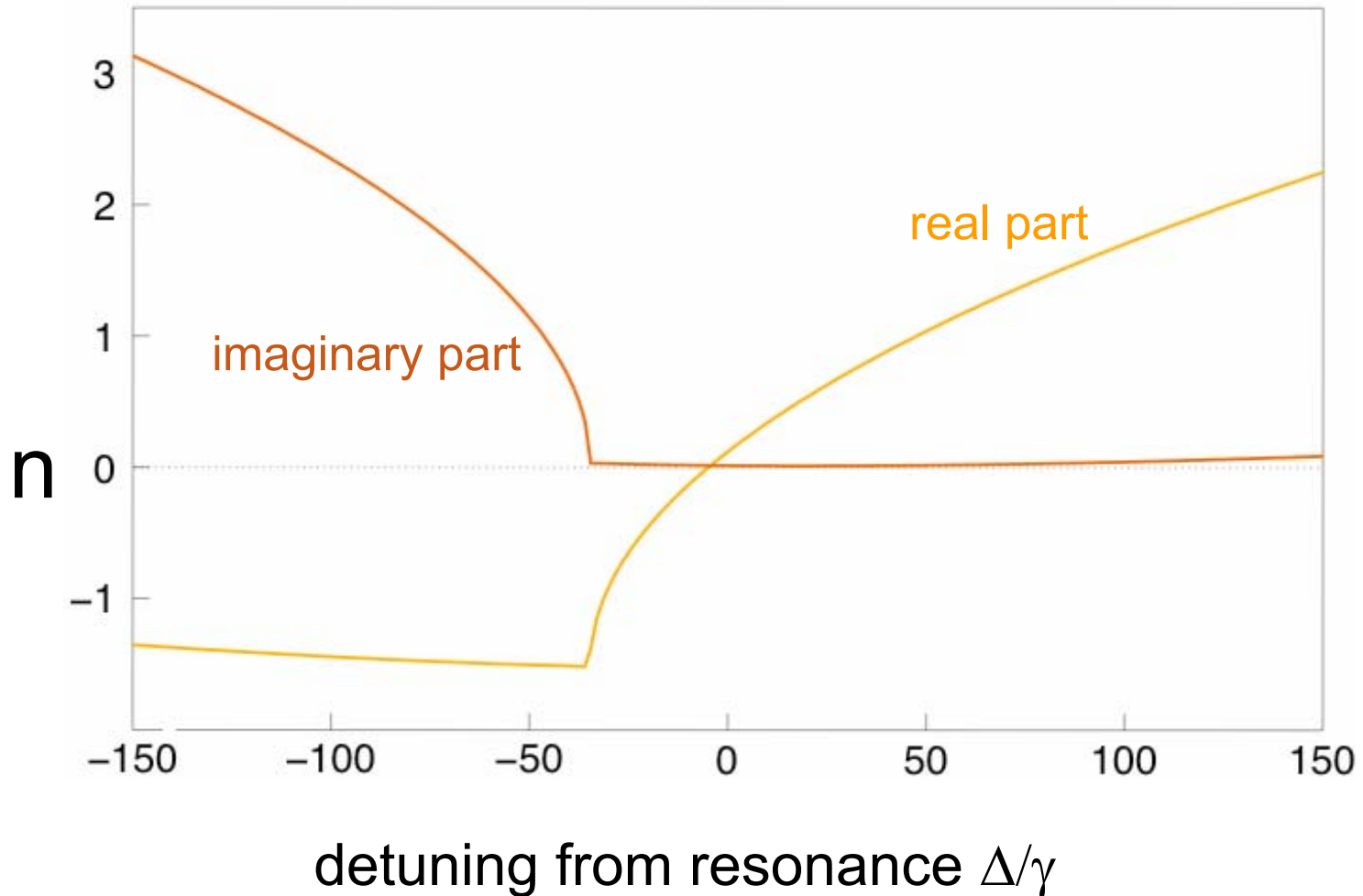


— real part

— imaginary part

Index of refraction

density $N = 5 \times 10^{16} \text{ cm}^{-3}$



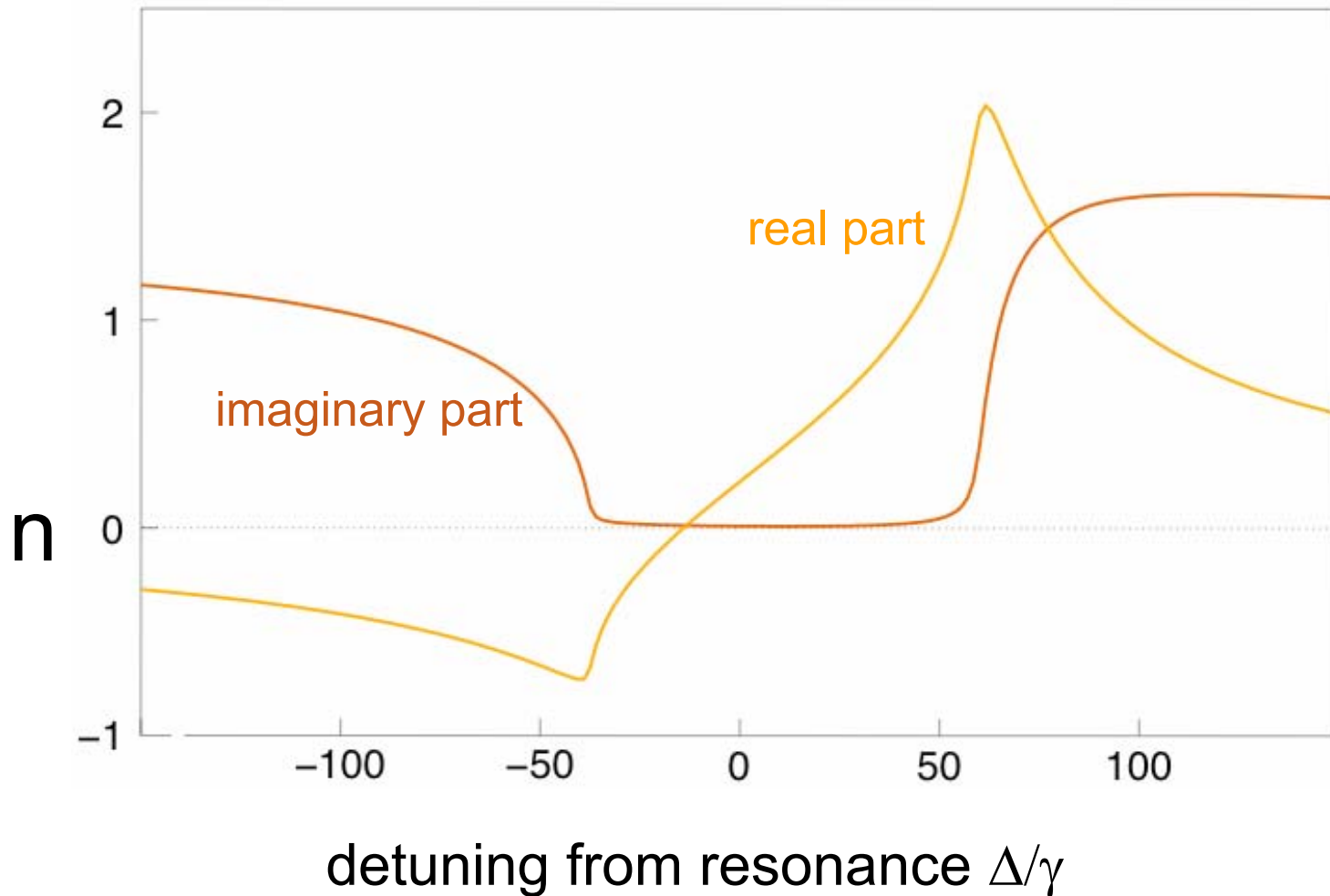
Local field corrections

$$\begin{aligned} E &\rightarrow E_{\text{loc}} = E + \frac{4\pi}{3}P \\ B &\rightarrow B_{\text{loc}} = B + \frac{4\pi}{3}M \end{aligned}$$

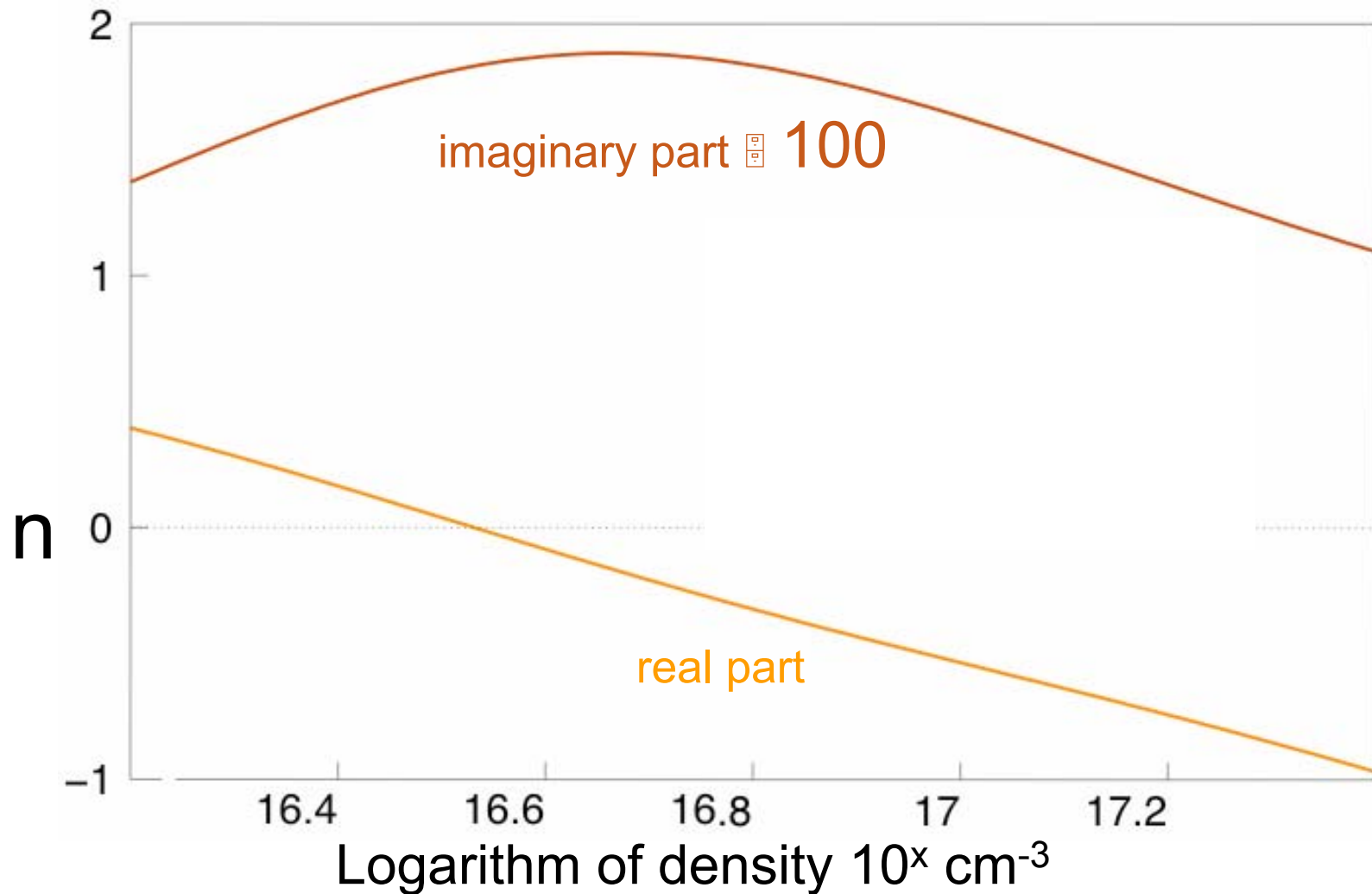
re-calculate χ 's and ξ 's . . .

Local field corrections

density $N = 5 \times 10^{16} \text{ cm}^{-3}$

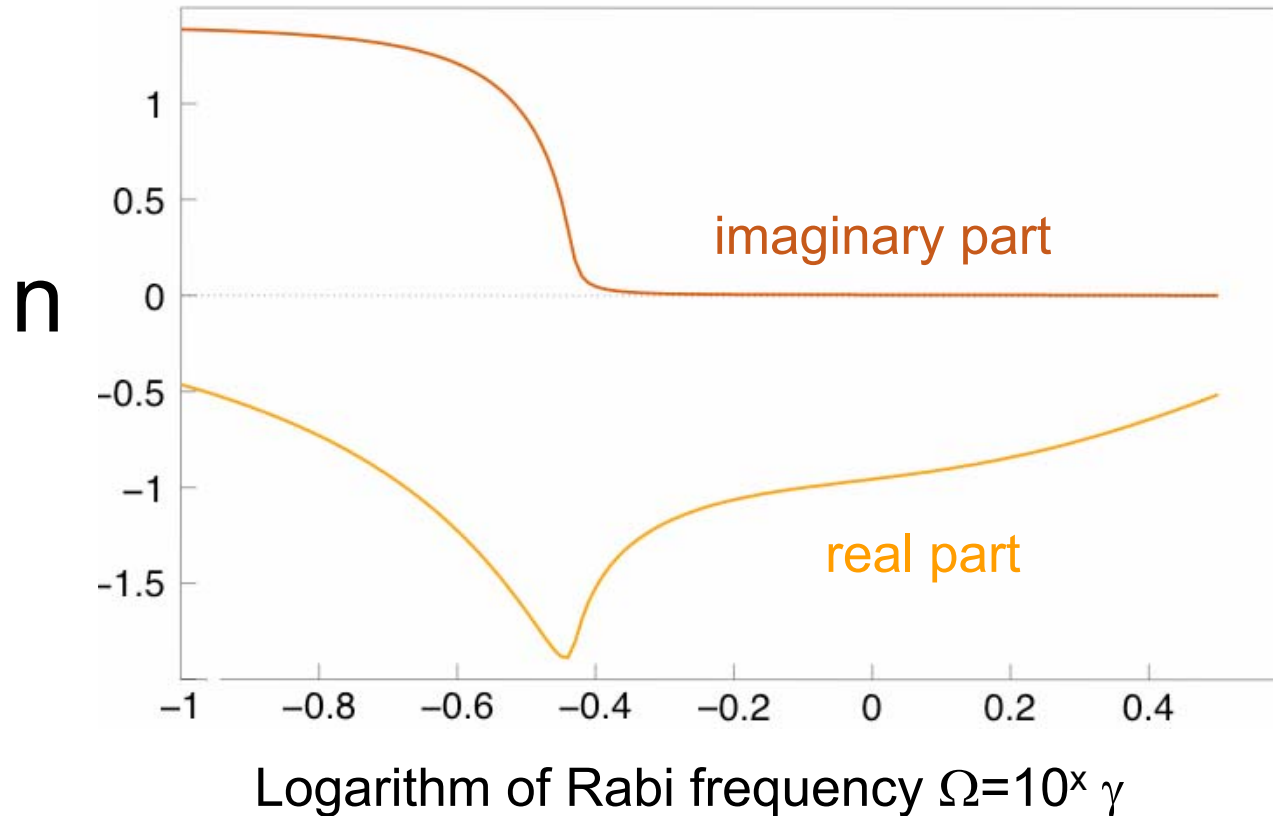


Density dependence



Fine tuning

n can be **tuned** by changing coupling field Rabi frequency Ω :



Application: e.g., for superlens, $n=-1$ is needed **exactly!**

Realization schemes

- **Atoms:** e.g. Neon
- **Molecules:** Use different rotational levels for different parities
- **Bound excitons:** use D^0 states with different parities for lower, and D^0X states with different parities for upper states.

Outlook

- **Materials:**
 - Problem of **high-frequency M1** transitions in atoms and molecules
 - **Parity** in solid state systems
- Dimension: **3D?**
- Comparison with “traditional” method + **gain**
- **Systems:**
 - Optimize level scheme
 - Utilize **tensorial** character of ε

Conclusions

- Use of negative refraction:
 - superlenses and others
- Metamaterials:
 - chiral media for presence of cross coupling
 - EIT for suppression of absorption
 - energy and Rabi freq. of coupling fields for tuning

Effects

normal refraction

phase velocity $v \approx c$

Group velocity $v_{gr} < v$

negative refraction

phase velocity $v \approx -c$

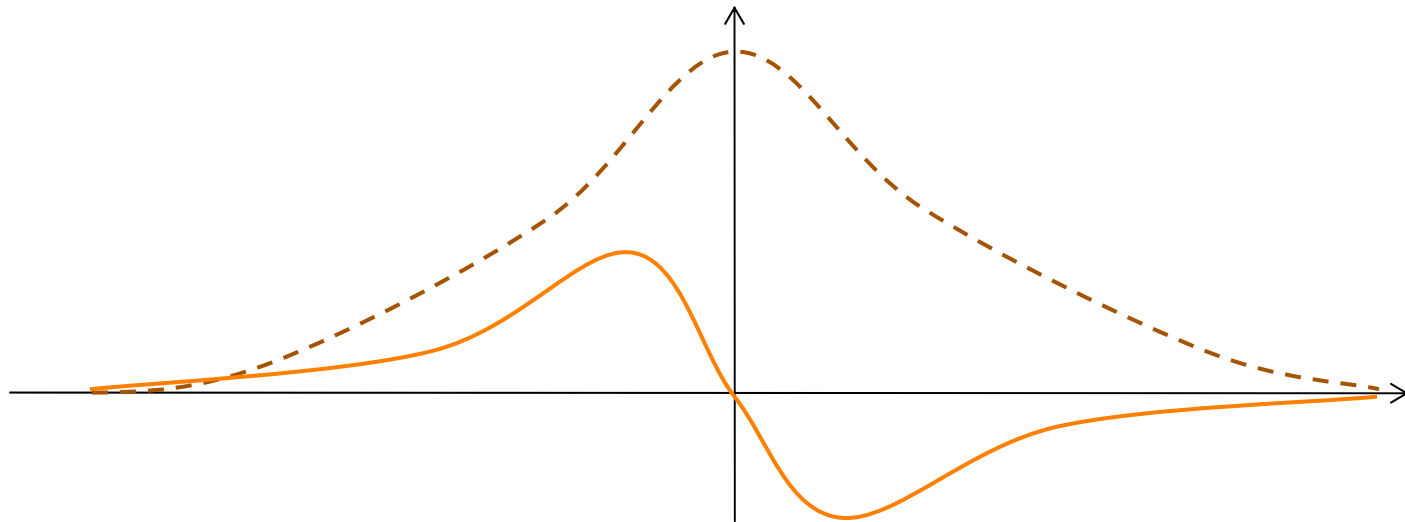
group velocity $v_{gr} \approx +c$

Problem: absorption

Kramers - Kronig:

relationship between refraction/absorption

large χ_e' (refraction) large χ_e'' (absorption)



Cross couplings

atomic picture:

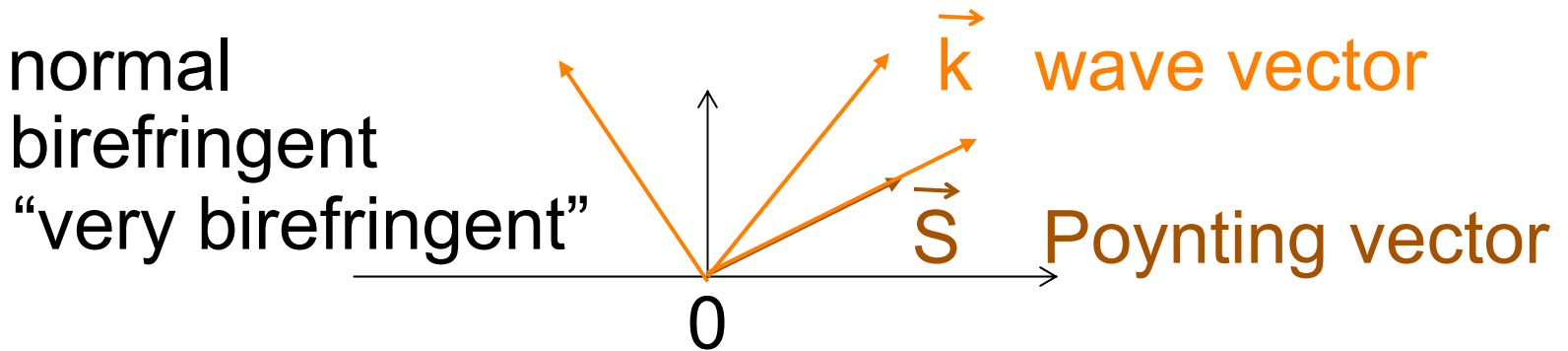
$$\rho_{34} = \alpha_{ee} \mathbf{E} + \alpha_{eb} \mathbf{B}$$

$$\rho_{21} = \alpha_{be} \mathbf{E} + \alpha_{bb} \mathbf{B}$$

Solve for α ...

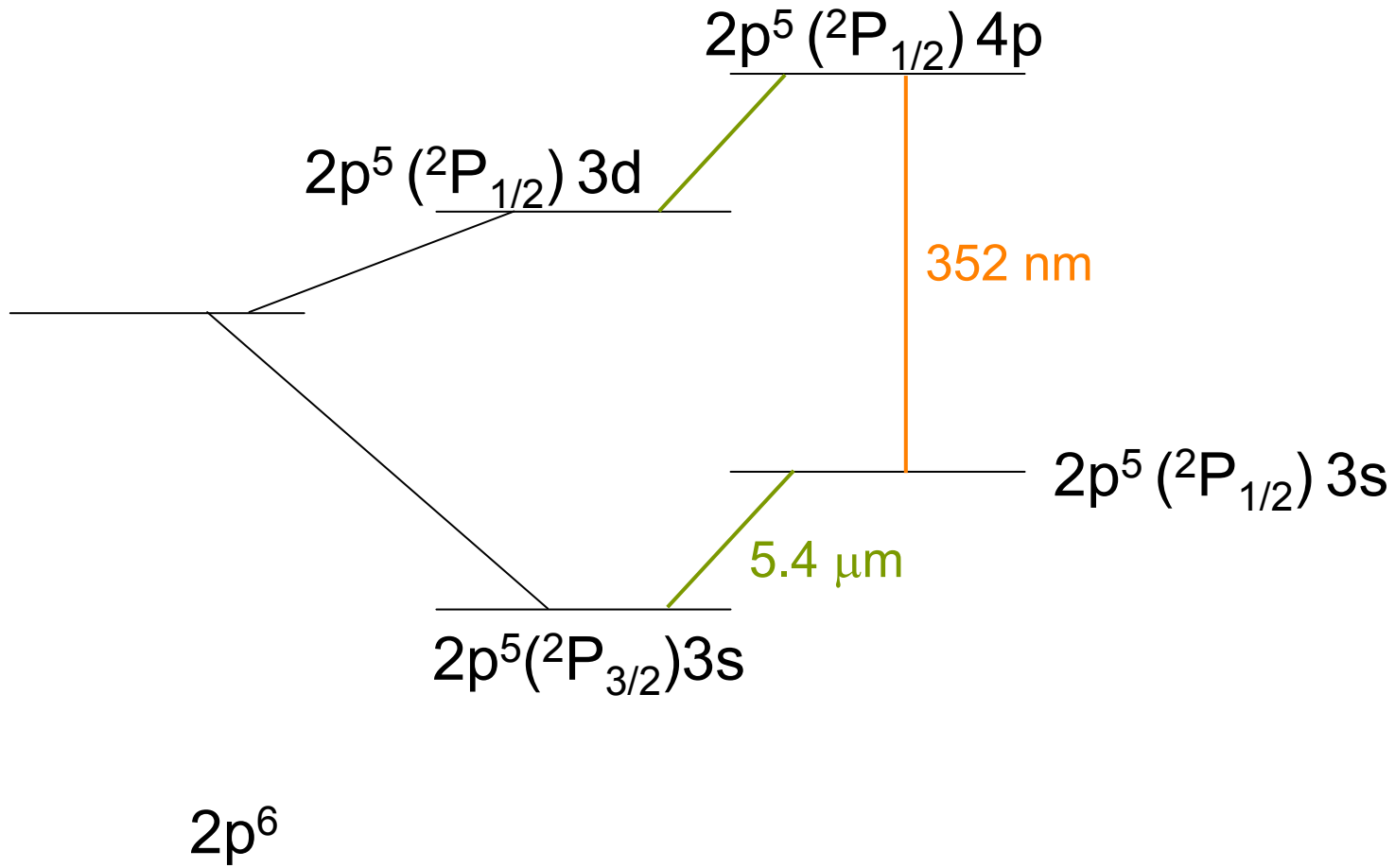
Different approach

- usual problem: μ (χ_m)
- Instead: leave μ and make ε into tensor (“geometric approach”)

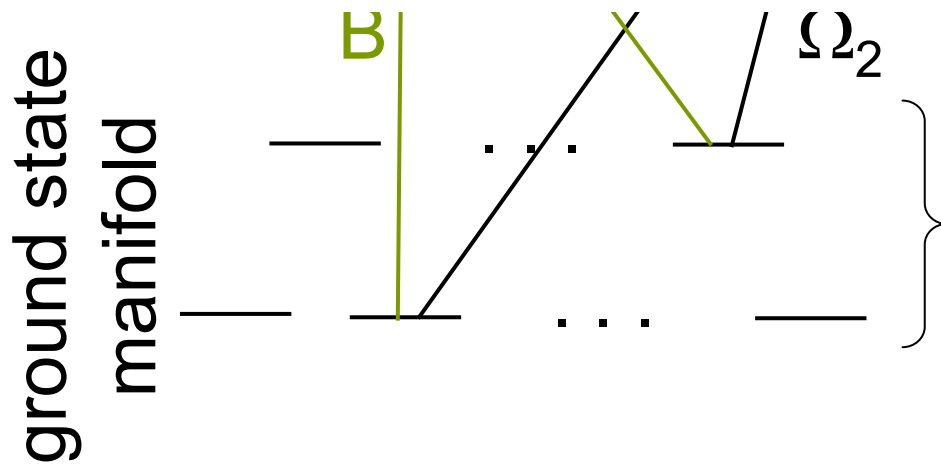


Disadvantage: works only in waveguide (i.e. 1D)

Neon



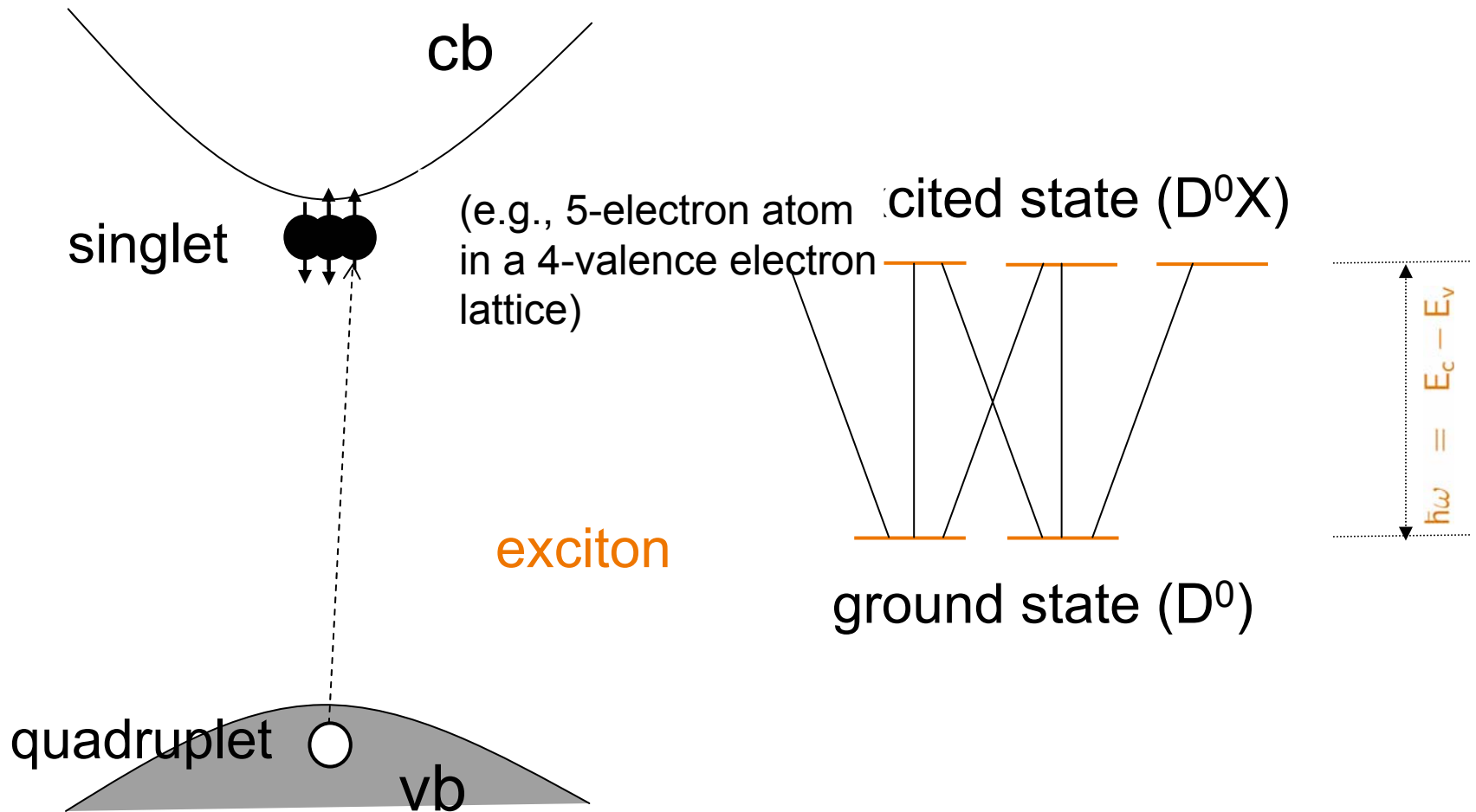
Molecular or solid state levels



one even, one odd parity (e.g., even and odd rotational level) for $|1 \textcircled{5}$ and $|4 \textcircled{5}$

Bound exciton

momentum picture:



Bound exciton

momentum picture:

